

lectures on
**DIFFERENTIAL GEOMETRY
 AND TOPOLOGY**

V2

Prof. M. Spina, UCSC Brescia

**LIE DERIVATIVES
 (FORM & GENERAL TENSOR CASE)**

Lecture XXIII

- Lie derivative for k -forms.
- Cartan's magic formula

Remark. Using a partition of unity, a form defined on $U \subset M$ (open) can be extended to a form defined on M .

We wish to give an intrinsic formulation of $d\omega$, $\omega \in \Lambda^1(M)$

$$d\omega(X, Y) = X\omega(Y) - Y\omega(X) - \omega([X, Y])$$

Proof. We work locally, so it is enough to take $\omega = u dv$, $u, v \in \mathcal{C}^\infty(M)$.

l.h.s. $d\omega = du \wedge dv$ and

$$d\omega(X, Y) = (du \wedge dv)(X, Y) = (du)(X)(dv)(Y) - (du)(Y)(dv)(X) \\ = X(u)Y(v) - X(v)Y(u)$$

r.h.s. $X\omega(Y) = X[u dv(Y)] = X[u Y(v)] \\ = X(u)Y(v) + u X(Y(v))$

$$-Y\omega(X) = -Y[u dv(X)] = -Y(u X(v)) \\ = -Y(u)X(v) - u Y(X(v))$$

they add to $u \cdot [X, Y](v)$

$$-\omega([X, Y]) = -u dv([X, Y]) = -u [X, Y](v)$$

Thus r.h.s. = ... l.h.s

In general, one shows that, for $\omega \in \Delta^k(M)$

$$d\omega(x_1 \dots x_{k+1}) = \sum_{i=1}^{k+1} (-1)^{i+1} x_i \omega(x_1 \dots \hat{x}_i \dots x_{k+1}) \\ + \sum_{i < j} (-1)^{i+j} \omega([x_i, x_j], x_1 \dots \hat{x}_i \dots \hat{x}_j \dots x_{k+1})$$

\uparrow
 first argument
 $\hat{} : \text{omission}$

★ Interior multiplication (of contraction)

$$i_X : \Delta^k(M) \rightarrow \Delta^{k-1}(M)$$

extended to $i : \Lambda \rightarrow \Lambda$

$$\omega \mapsto i_X \omega$$

$X \in \mathfrak{X}(M)$

$$(i_X \omega)(x_1 \dots x_{k-1})$$

actually, i is a purely algebraic operation.

$$:= \omega(\overset{\uparrow}{\cancel{X}}, x_1, \dots, x_{k-1})$$

\uparrow
 first slot

(other notation: $X \lrcorner \omega$)

i_X is linear and, for $X, Y \in \mathfrak{X}(M)$, $\alpha, \beta \in \mathbb{R}$

$$i_{\alpha X + \beta Y} \omega = \alpha i_X \omega + \beta i_Y \omega$$

$$(\diamond) \quad \boxed{i_X^2 = 0}$$

and i is an antiderivation

$(\diamond\diamond)$

$$\boxed{i_X(\omega \wedge \eta) = i_X \omega \wedge \eta + (-1)^k \omega \wedge i_X \eta}$$

\uparrow
 Δ^k

i behaves like d

Let us check that $\partial_x^2 = 0$ $w \in \Delta^{\mathbb{R}}$

$$\begin{aligned} \partial_x (\partial_x w) (x_1, \dots, x_{k-2}) &= \\ &= \partial_x w (x, x_1, \dots, x_{k-2}) = w(x, x, x_1, \dots, x_{k-2}) = 0 \end{aligned}$$

As for (ii), it is enough to check the formula

$$\begin{aligned} \partial_x (w^1 \wedge \dots \wedge w^k) &= \partial_x w^1 \wedge w^2 \wedge \dots \wedge w^k \\ &\quad - w^1 \wedge \partial_x w^2 \wedge \dots \wedge w^k \\ &\quad + w^1 \wedge w^2 \wedge \dots \wedge \partial_x w^k \end{aligned}$$

↑ ↑
1-forms

← quite useful in explicit computations

and this is true by Laplace's formula:

if $x_1 = x$.

$$\begin{aligned} (w^1 \wedge w^2 \wedge \dots \wedge w^k) (x_1, x_2, \dots, x_k) &= \det (w^i(x_j)) \\ &= \sum_{i=1}^k (-1)^{i-1} w^i(x_1) \underbrace{(w^1 \wedge \dots \wedge \hat{w}^i \wedge \dots \wedge w^k)}_{\det \begin{matrix} \times & & \\ & \times & \\ & & \times \end{matrix}} (x_2, \dots, x_k) \end{aligned}$$

* Lie derivative of a \mathbb{R} -form

Let $X \in \mathfrak{X}(M)$, $\omega \in \Lambda^k(M)$

The Lie derivative of ω along X , $L_X \omega \in \Lambda^k(M)$ is defined as:

recall that the pull-back is defined algebraically

F_t^X : flow of X

$$(F_t^X)^* \omega(F_t^X(P))$$

$$(L_X \omega)(P) = \left. \frac{d}{dt} [(F_t^X)^* \omega] \right|_{t=0} = \lim_{t \rightarrow 0} \frac{[(F_t^X)^* \omega](P) - \omega(P)}{t}$$

These two vectors can be compared

One may prove that $L_X(\omega \wedge \tau) =$

$$L_X \omega \wedge \tau + \omega \wedge L_X \tau$$

and, more generally

$$L_X(T \otimes S) = L_X T \otimes S + T \otimes L_X S \quad (*)$$

for any tensor fields

(generalized Leibniz rule)

also; for a $(0, k)$ -tensor one has:

$$(L_X \sigma)(Y_1, \dots, Y_k) = X(\sigma(Y_1, \dots, Y_k)) - \sigma([X, Y_1], Y_2, \dots, Y_k) - \dots - \sigma(Y_1, \dots, Y_{k-1}, [X, Y_k])$$

Thus, in turn, coming from

$$X(\sigma(Y_1, \dots, Y_k)) = (L_X \sigma)(Y_1, \dots, Y_k) + \sigma(L_X Y_1, Y_2, \dots, Y_k) + \dots + \sigma(Y_1, \dots, L_X Y_k)$$

$L_X(\sigma(Y_1, \dots, Y_k))$ Leibniz rule: all arguments are varied one at a time

Let us check (*) in the case of a tensor product of covariant tensors, just to pinpoint the (quite simple) basic idea.

$$\mathcal{L}_x (\sigma \otimes \tau)(P) = \lim_{t \rightarrow 0} \frac{(F_t^x)^* [(\sigma \otimes \tau)(F_t^x(P))] - (\sigma \otimes \tau)(P)}{t}$$

$$= \lim_{t \rightarrow 0} \frac{\overbrace{(F_t^x)^* \sigma(F_t^x(P)) \otimes (F_t^x)^* \tau(F_t^x(P))}^A - \underbrace{\sigma(P) \otimes \tau(P)}_B}{t}$$

$$= \lim_{t \rightarrow 0} \left[\frac{\overbrace{A - \sigma(P) \otimes (F_t^x)^* \tau(F_t^x(P))}^C}{t} + \frac{C - B}{t} \right]$$

[This is the same idea one uses to prove $(fg)' = f'g + fg'$ in calculus]

$$= \lim_{t \rightarrow 0} \frac{\overbrace{(F_t^x)^* \sigma(F_t^x(P)) - \sigma(P)}^{\mathcal{L}_x \sigma} \otimes \overbrace{(F_t^x)^* \tau(F_t^x(P))}^{\tau(P)}}{t} + \sigma(P) \otimes \lim_{t \rightarrow 0} \frac{\overbrace{(F_t^x)^* \tau(F_t^x(P)) - \tau(P)}^{\mathcal{L}_x \tau}}{t}$$

$$= \mathcal{L}_x \sigma \otimes \tau + \sigma \otimes \mathcal{L}_x \tau$$

$f(x+h)g(x+h) - f(x)g(x)$
 $= f(x+h)g(x+h) - f(x)g(x+h) + f(x)g(x+h) - f(x)g(x)$
 $= \{f(x+h) - f(x)\}g(x+h) + f(x)\{g(x+h) - g(x)\}$
 \vdots

also recall that in general if $S = \sigma_J^I$, $T = \tau_L^K$

$$S \otimes T = \sigma_J^I \tau_L^K \leftarrow \text{no summation, just take all products}$$

↑
numerical product

* Theorem Let $\omega \in \Delta(M)$, $X \in \mathcal{X}(M)$.

One has

$$\boxed{L_X \omega = d i_X \omega + i_X d\omega}$$

*** (Cartan's magic formula)

⚠ It holds for forms, not for general tensors (d is not defined...)

Pf. We prove it for $\omega = v dv$. The general case follows by induction.

First of all, notice that

$$(L_X dv)(Y) = X [dv(Y)] - dv([X, Y])$$

general
formula

$$= X[Y(v)] - [X, Y](v) = (XY - YX)(v) = Y X(v) - L_X v$$

$$= d(L_X v)(Y)$$

namely

$$\boxed{L_X dv = d L_X v}$$

As a corollary, if (x^1, \dots, x^n) are local coordinates and $X = \frac{\partial}{\partial x^1}$,

just to fix ideas, then $L_X(dx^1, \dots, dx^k) = 0$,

This following from $L_X dx^j = d L_X x^j = d \left(\frac{\partial x^j}{\partial x^1} \right)$

and from the general Leibniz rule.

$$= d(\delta_{1j}) = 0$$

Now compute:

$$\begin{aligned}
 L_X(u dv)(Y) &= (L_X u \cdot dv + u L_X dv)(Y) \\
 &= (L_X u dv + u d L_X v)(Y) \\
 &= (X(u) dv + u d(X(v)))(Y) \\
 &= X(u) Y(v) + u Y(X(v)) \quad \textcircled{I}
 \end{aligned}$$

compute $\textcircled{II} =$

$$\begin{aligned}
 &= (d i_X + i_X d)(u dv)(Y) = \\
 &= d(i_X(u dv))(Y) + i_X d(u dv)(Y) \\
 &= d(\underbrace{i_X u}_{=0} dv + u \underbrace{i_X dv}_{X(v)})(Y) + i_X(\underbrace{du \wedge dv}_{(du \wedge dv)(X, Y)})(Y) \\
 &= d(u X(v))(Y) + X(u) Y(v) - X(v) Y(u) \\
 &= (du X(v) + u d(X(v)))(Y) + X(u) Y(v) - X(v) Y(u) \\
 &= Y(u) X(v) + u Y(X(v)) + X(u) Y(v) - X(v) Y(u) \\
 &= X(u) Y(v) + u Y(X(v)) = \textcircled{I} \quad \square
 \end{aligned}$$

Notice that, in general, for forms, one has:

$$\boxed{L_X d = d L_X}$$

Indeed: $L_X d = (d i_X + i_X d) d = d i_X d + i_X d^2 = d i_X d$

whereas $d L_X = d(d i_X + i_X d) = \underbrace{d^2 i_X}_{=0} + d i_X d = d i_X d \quad \square$

Examples

In \mathbb{R}^3

1. $X = \frac{z}{\partial z}$ $\omega = a \, dx \wedge dy$

$a \in C^1(\mathbb{R}^3)$

$$d\omega = da \wedge dx \wedge dy = \frac{\partial a}{\partial z} dz \wedge dx \wedge dy = \frac{\partial a}{\partial z} dx \wedge dy \wedge dz$$

$$\left[i_X d\omega = \frac{\partial a}{\partial z} dy \wedge dz \right]$$

$$(i_X d\omega)(y, z) = \left(\frac{\partial a}{\partial z} dx \wedge dy \wedge dz \right) (x, y, z) =$$

↖ R running through $(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z})$

$$= \frac{\partial a}{\partial z} \begin{vmatrix} \overset{1}{dx(x)} & dx(y) & dx(z) \\ dy(x) & & \\ dz(x) & & \end{vmatrix} = \frac{\partial a}{\partial z} \begin{vmatrix} dy(y) & dy(z) \\ dz(y) & dz(z) \end{vmatrix}$$

$$= \frac{\partial a}{\partial z} (dy \wedge dz)(y, z)$$

One can proceed directly from the definition

$$\frac{\partial a}{\partial z} (dx \wedge dy \wedge dz)(x, \bullet, \blacklozenge) =$$

$$= \frac{\partial a}{\partial z} \left(\overset{1}{dx(x)} dy(\bullet) dz(\blacklozenge) - \underbrace{dx(x)}_1 dy(\blacklozenge)(\bullet) + \dots \right)$$

all other terms are 0
since $dy(x) = dz(x) = 0$

$$= \frac{\partial a}{\partial z} dy \wedge dz(\bullet, \blacklozenge)$$

* Actually, the one below is the quickest method:

$$i_X \left(\frac{\partial a}{\partial z} dx \wedge dy \wedge dz \right) = \frac{\partial a}{\partial z} i_X (dx \wedge dy \wedge dz)$$

$$= \frac{\partial a}{\partial z} (i_X dx) \wedge dy \wedge dz - dx \wedge i_X dy \wedge dz + \dots = 0$$

$dy(x) = dx(y) = \frac{\partial y}{\partial x} = 0$

$$= \frac{\partial a}{\partial z} dy \wedge dz$$

$$i_X \omega = a \underbrace{i_X}_{\frac{\partial}{\partial x}}(dx \wedge dy) = a dy$$

(proceed as before..)

$$\boxed{d(i_X \omega) = d(a dy) = da \wedge dy = \frac{\partial a}{\partial x} dx \wedge dy + \frac{\partial a}{\partial z} dz \wedge dy}$$

$$= \left[\frac{\partial a}{\partial x} dx \wedge dy - \frac{\partial a}{\partial z} dy \wedge dz \right]$$

Therefore $(d i_X + i_X d) \omega = \frac{\partial a}{\partial x} dx \wedge dy$

Compute $L_X \omega$ directly:

$$L_X (a dx \wedge dy) = \underbrace{(L_X a)}_{\substack{x(a) \\ = \frac{\partial a}{\partial x}}} dx \wedge dy + a \underbrace{L_X dx \wedge dy}_{=0}$$

$$= \frac{\partial a}{\partial x} dx \wedge dy \quad \checkmark$$

Leibniz

2. $X = \frac{\partial}{\partial x} \quad \omega = a dy \wedge dz$

$$d\omega = da \wedge dy \wedge dz = \frac{\partial a}{\partial x} dx \wedge dy \wedge dz$$

$$i_X d\omega = \dots = \frac{\partial a}{\partial x} dy \wedge dz$$

$$i_X \omega = \dots = 0 \quad (a dy \wedge dz)(X, \cdot) = a [dy(X) dz(\cdot) - dz(\cdot) dy(X)]$$

$$d i_X \omega = 0 \quad (d i_X + i_X d) \omega = \frac{\partial a}{\partial x} dx \wedge dy \wedge dz$$

$$L_X (a dy \wedge dz) = L_X a \cdot dy \wedge dz + \dots = \frac{\partial a}{\partial x} dy \wedge dz \quad \checkmark$$

$3 \cdot$ in \mathbb{R}^3 $\omega = dx \wedge dy \wedge dz$ standard volume form

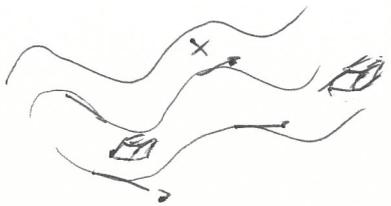
Let $X = \alpha \frac{\partial}{\partial x} + \beta \frac{\partial}{\partial y} + \gamma \frac{\partial}{\partial z}$ $\alpha, \beta, \gamma \in C^\infty(\mathbb{R}^3)$

Compute $L_X \omega = \left(\frac{\partial \alpha}{\partial x} + \frac{\partial \beta}{\partial y} + \frac{\partial \gamma}{\partial z} \right) dx \wedge dy \wedge dz$

(Computes variation of volume elements)

divergence of X

 metric is involved



$L_X \omega = 0$ if $\text{div } X = 0$ X solenoidal (or divergence-free)

defines volume-preserving flows

\equiv incompressible flows

This holds on Riemannian manifolds

Let us perform the calculation directly:

$$L_X(dx \wedge dy \wedge dz) = L_X dx \wedge dy \wedge dz + dx \wedge L_X dy \wedge dz + dx \wedge dy \wedge L_X dz$$

$$= d(L_X x) \wedge dy \wedge dz + \text{similar terms}$$

$$= d(X(x)) \wedge dy \wedge dz + \dots$$

$$X(x) = \dots = \alpha$$

$$= d\alpha \wedge dy \wedge dz + \dots$$

$$= \frac{\partial \alpha}{\partial x} dx \wedge dy \wedge dz + \dots \quad \checkmark$$

Use Cartan

$$d\omega = 0 \Rightarrow L_X \omega = d \iota_X \omega$$

a 4-form on \mathbb{R}^3

$$\iota_{X_1} \omega = \dots = \alpha dy \wedge dz \Rightarrow d \iota_{X_1} \omega = \frac{\partial \alpha}{\partial x} dx \wedge dy \wedge dz$$

"
 $\alpha \frac{\partial}{\partial x}$

Summing up, we achieve the conclusion. \checkmark

4. Let (M, g) be a Riemannian manifold and $X \in \mathcal{X}(M)$. Compute $L_X g$. *Work locally:*

$$X = \xi^i \frac{\partial}{\partial x^i}$$

(we cannot use Cartan's formula, g is not a 2-form; in fact, it is a (symmetric, positive definite) $(0,2)$ -tensor field)

$$L_X (g_{ij} dx^i dx^j) =$$

$$= L_X (g_{ij}) dx^i dx^j + g_{ij} (L_X dx^i) dx^j + g_{ij} dx^i L_X dx^j$$

\parallel \parallel \parallel
 $X(g_{ij})$ $dL_X x^i$ $d\xi^j$
 \parallel \parallel \parallel
 $dL_X x^i$ $d\xi^j$
 \parallel \parallel
 $dL_X x^i$ $d\xi^j$
 \parallel \parallel
 $dL_X x^i$ $d\xi^j$

$$= X(g_{ij}) dx^i dx^j + g_{ij} d\xi^i dx^j + g_{ij} dx^i d\xi^j$$

$$= \sum^{\mathbb{R}} \frac{\partial g_{ij}}{\partial x^k} dx^i dx^j + g_{ij} \frac{\partial \xi^i}{\partial x^k} dx^k dx^j + g_{ij} dx^i \frac{\partial \xi^j}{\partial x^k} dx^k$$

one wants ij
*so relabel $\mathbb{R} \rightarrow i$
 $i \rightarrow k$*

$$= \sum^{\mathbb{R}} \frac{\partial g_{ij}}{\partial x^k} dx^i dx^j + g_{kj} \frac{\partial \xi^k}{\partial x^i} dx^i dx^j + g_{ik} \frac{\partial \xi^k}{\partial x^j} dx^i dx^j$$

*relabel: $k \rightarrow j$
 $j \rightarrow k$*

$$= \left(\sum^{\mathbb{R}} \frac{\partial g_{ij}}{\partial x^k} + g_{kj} \frac{\partial \xi^k}{\partial x^i} + g_{ik} \frac{\partial \xi^k}{\partial x^j} \right) dx^i dx^j$$

X is a Killing vector field (or, simply, X is Killing)

if $L_X g = 0$ ($\Rightarrow X$ generates isometries of (M, g))

$$\sum^{\mathbb{R}} \frac{\partial g_{ij}}{\partial x^k} + g_{kj} \frac{\partial \xi^k}{\partial x^i} + g_{ik} \frac{\partial \xi^k}{\partial x^j} = 0$$

i, j fixed
summation
over \mathbb{R}

For instance, in $(\mathbb{R}^n, g = \sum (dx^i)^2)$

$$g_{ij} = \delta_{ij}$$

The Killing condition becomes:

$$\left(\delta_{kj} \frac{\partial \xi^k}{\partial x^i} + \delta_{ik} \frac{\partial \xi^k}{\partial x^j} \right) = 0$$

i.e.

$$\boxed{\frac{\partial \xi^j}{\partial x^i} + \frac{\partial \xi^i}{\partial x^j} = 0}$$

in particular
 $\frac{\partial \xi^i}{\partial x^i} = 0$

Constant vector fields and generators of rotations in

the plane (i, j) , namely $x^i \frac{\partial}{\partial x^j} - x^j \frac{\partial}{\partial x^i}$,

$i \neq j$

i, j fixed

are examples of Killing fields:

$$\xi^i = -x^j \quad \xi^j = x^i$$

$$\frac{\partial \xi^j}{\partial x^i} + \frac{\partial \xi^i}{\partial x^j} = \frac{\partial x^j}{\partial x^i} - \frac{\partial x^i}{\partial x^j} = 1 - 1 = 0$$