

# CONCENTRATING SOLUTIONS FOR FRACTIONAL SCHRÖDINGER-POISSON SYSTEMS WITH CRITICAL GROWTH

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ABSTRACT. We consider a class of fractional Schrödinger-Poisson systems with critical growth

$$\begin{cases} (-\Delta)^s u + \lambda V(x)u + \phi u = f(u) + |u|^{2_s^* - 2} u, & x \in \mathbb{R}^3, \\ (-\Delta)^t \phi = u^2, & x \in \mathbb{R}^3, \end{cases}$$

where  $s, t \in (0, 1)$  with  $2s + 2t > 3$ ,  $\lambda > 0$  denotes a parameter,  $V : \mathbb{R}^3 \rightarrow \mathbb{R}$  admits a potential well  $\Omega \triangleq \text{int}V^{-1}(0)$  and  $2_s^* \triangleq \frac{6}{3-2s}$  is the fractional Sobolev critical exponent. Under certain assumptions on  $f$ , we obtain the existence and concentrating behavior of nontrivial solutions using variational methods.

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## 1. INTRODUCTION

**1.1. Overview.** In this article, we investigate the existence and concentration of nontrivial solutions for the following fractional Schrödinger-Poisson system with critical growth

$$(1.1) \quad \begin{cases} (-\Delta)^s u + \lambda V(x)u + \phi u = f(u) + |u|^{2_s^* - 2} u, & x \in \mathbb{R}^3, \\ (-\Delta)^t \phi = u^2, & x \in \mathbb{R}^3, \end{cases}$$

where  $s, t \in (0, 1)$  with  $2s + 2t > 3$ ,  $\lambda > 0$  denotes a parameter,  $V : \mathbb{R}^3 \rightarrow \mathbb{R}$  admits a potential well  $\Omega \triangleq \text{int}V^{-1}(0)$  and  $2_s^* \triangleq \frac{6}{3-2s}$  is the fractional Sobolev critical exponent. On the potential  $V$ , we shall firstly make the following assumptions

- (V<sub>1</sub>)  $V \in \mathcal{C}(\mathbb{R}^3, \mathbb{R})$  with  $V \geq 0$  on  $\mathbb{R}^3$ ;
- (V<sub>2</sub>) there is  $c > 0$  such that the set  $\Sigma \triangleq \{x \in \mathbb{R}^3 : V(x) < c\}$  has positive finite Lebesgue measure;
- (V<sub>3</sub>)  $\Omega = \text{int}V^{-1}(0)$  is nonempty with smooth boundary with  $\bar{\Omega} = V^{-1}(0)$ ,  $V^{-1}(0) \triangleq \{x : V(x) = 0\}$ .

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In celebrated papers, Bartsch and his collaborators initially proposed the above hypotheses to study the nonlinear Schrödinger equations, see [5, 6]. As we all know, the harmonic trapping potential

$$V(x) = \begin{cases} \omega_1|x_1|^2 + \omega_2|x_2|^2 + \omega_3|x_3|^2 - \omega, & \text{if } |(\sqrt{\omega_1}x_1, \sqrt{\omega_2}x_2, \sqrt{\omega_3}x_3)|^2 \geq \omega, \\ 0, & \text{if } |(\sqrt{\omega_1}x_1, \sqrt{\omega_2}x_2, \sqrt{\omega_3}x_3)|^2 \leq \omega, \end{cases}$$

with  $\omega > 0$  satisfying  $(V_1)$ - $(V_3)$ , where  $\omega_i > 0$  is called by the anisotropy factor of the trap in quantum physics and trapping frequency of the  $i$ th-direction in mathematics, see e.g. [7, 12, 24]. Indeed, the potential  $\lambda V$ , instead of  $V$ , with assumptions  $(V_1)$ - $(V_3)$  can be read as the steep potential well.

Over the past several decades, there was considerable attention to the standing, or solitary, wave solutions of Schrödinger-Poisson systems of the type

$$(1.2) \quad \begin{cases} i \frac{\partial \psi}{\partial t} = \Delta \psi - W(x)\psi + \phi \psi + \tilde{g}(|\psi|)\psi, & \text{in } \mathbb{R}^+ \times \mathbb{R}^3, \\ -\Delta \phi = |\psi|^2, & \text{in } \mathbb{R}^3, \end{cases}$$

where  $\psi : \mathbb{R}^3 \times \mathbb{R} \rightarrow \mathbb{C}$  is the time-dependent wave function,  $W : \mathbb{R}^3 \rightarrow \mathbb{R}$  stands for the real external potential,  $\phi$  represents an internal potential for a nonlocal self-interaction of wave function and nonlinear term  $g(\psi) \triangleq \tilde{g}(|\psi|)\psi$  describes the interaction effect among particles. By inserting the standing wave ansatz  $\psi(x, t) = \exp(-i\omega t)u(x)$  with  $\omega \in \mathbb{R}$  and  $x \in \mathbb{R}^3$  into (1.2), then  $u : \mathbb{R}^3 \rightarrow \mathbb{R}$  satisfies the Schrödinger-Poisson system

$$(1.3) \quad \begin{cases} -\Delta u + \bar{W}(x)u + \phi u = g(u), & \text{in } \mathbb{R}^3, \\ -\Delta \phi = u^2, & \text{in } \mathbb{R}^3, \end{cases}$$

where and in the sequel  $\bar{W}(x) = W(x) + \omega$  for all  $x \in \mathbb{R}^3$ . We refer the interested readers to [8, 9] and the references therein for more about the physical background of (1.2). There are many interesting works about the existence of positive solutions, positive ground states, multiple solutions, sign-changing solutions and semiclassical states to system (1.3), see e.g. [2, 4, 15, 16, 28, 29, 38, 46] and their references therein.

In [18], Jiang and Zhou firstly applied the steep potential well to the Schrödinger-Poisson system and proved the existence of nontrivial solutions and ground state solutions. Subsequently, by using the linking theorem [27, 42], the authors in [45] considered the existence and concentration of nontrivial solutions for the following Schrödinger-Poisson system

$$(1.4) \quad \begin{cases} -\Delta u + \lambda V(x)u + K(x)\phi u = |u|^{p-2}u, & x \in \mathbb{R}^3, \\ -\Delta \phi = K(x)u^2, & x \in \mathbb{R}^3, \end{cases}$$

under the following conditions:

( $\tilde{V}$ )  $V \in C(\mathbb{R}^3, \mathbb{R})$  and  $V$  is bounded from below;

and  $(V_2)$ - $(V_3)$  with some suitable assumptions on  $K : \mathbb{R}^3 \rightarrow \mathbb{R}$  for  $4 \leq p < 6$ . It is worth mentioning that, in particular, they investigated the existence and concentration of nontrivial solutions to (1.4) by the monotonicity trick due to Jeanjean [17] under the conditions  $(V_1)$ - $(V_3)$ ,  $K \in L^\infty_{\text{loc}}(\mathbb{R}^3) \cap L^2(\mathbb{R}^3)$  and

( $\bar{V}$ )  $V(x)$  is weakly differentiable such that  $(x, \nabla V) \in L^{p_1}(\mathbb{R}^3)$  for some  $p_1 \in [\frac{3}{2}, \infty]$ , and

$$2V(x) + (x, \nabla V) \geq 0 \text{ for a.e. } x \in \mathbb{R}^3,$$

where  $(\cdot, \cdot)$  is the usual inner product in  $\mathbb{R}^3$ .

( $\bar{K}$ )  $K(x)$  is weakly differentiable such that  $(x, \nabla K) \in L^{p_2}(\mathbb{R}^3)$  for some  $p_2 \in [2, \infty]$ , and

$$\frac{2(p-3)}{p}K(x) + (x, \nabla K) \geq 0 \text{ for a.e. } x \in \mathbb{R}^3.$$

Whereas, the related researches on fractional Schrödinger-Poisson systems like (1.1) are not as rich as the classic Schrödinger-Poisson system (1.3). Actually, we shall reach the system (1.1) by supposing  $s = t = 1$  and  $K(x) \equiv 1$  for each  $x \in \mathbb{R}^3$  in the system (1.4). As a consequence, one of the aims in this

paper is to generalize the corresponding results obtained in [45] to the fractional case which makes the studies interesting.

When it comes to the fractional order operators, the following fractional Schrödinger equation

$$(1.5) \quad (-\Delta)^\alpha u + V(x)u = f(x, u), \quad x \in \mathbb{R}^N,$$

is usually used to study the standing wave solutions  $\psi(x, t) = u(x)e^{-i\omega t}$  for the equation

$$i\hbar \frac{\partial \psi}{\partial t} = \hbar^2 (-\Delta)^\alpha \psi + W(x)\psi - f(x, \psi), \quad x \in \mathbb{R}^N,$$

where  $\hbar$  is the Planck's constant,  $W : \mathbb{R}^N \rightarrow \mathbb{R}$  is an external potential and  $f$  a suitable nonlinearity. Because the fractional Schrödinger equation appears in problems involving nonlinear optics, plasma physics and condensed matter physics, it is one of the main objects of the fractional quantum mechanic. The equation (1.5) has been firstly proposed by Laskin [20, 21] as a result of expanding the Feynman path integral, from the Brownian-like to the Lévy-like quantum mechanical paths. In [11], Caffarelli and Silvestre transformed the nonlocal operator  $(-\Delta)^\alpha$  to a Dirichlet-Neumann boundary value problem for a certain elliptic problem with local differential operators defined on the upper half space. This technique is a powerful tool to deal with the equations involving fractional operators in the respects of regularity and variational methods, please see [1, 15] and their references for example. When the conditions  $(V_1)$ - $(V_3)$  are satisfied, Yang and Liu [43] established the multiplicity and concentration of solutions for the following fractional Schrödinger equation

$$(-\Delta)^\alpha u + \lambda V(x)u = f(x, u) + g(x)|u|^{v-2}u, \quad x \in \mathbb{R}^N,$$

involving a  $k$ -order asymptotically linear term  $f(x, u)$ , where  $s \in (0, 1)$ ,  $2s < N$ ,  $1 \leq k < 2_s^* - 1 = \frac{N+2s}{N-2s}$  and  $g \in L^{\frac{v}{2-v}}(\mathbb{R}^N)$  with  $1 < v < 2$ . There exist some other meaningful results in [3, 10, 34] and their references on fractional Schrödinger equations.

Recently, Teng [39] contemplated the existence of ground state solutions to the following fractional Schrödinger-Poisson system

$$\begin{cases} (-\Delta)^s u + V(x)u + \phi u = |u|^{p-2}u + \mu|u|^{2_s^*-2}u, & x \in \mathbb{R}^3, \\ (-\Delta)^t \phi = u^2, & x \in \mathbb{R}^3, \end{cases}$$

where the potential  $V : \mathbb{R}^3 \rightarrow \mathbb{R}^+$  satisfies some technical assumptions,  $\mu = 1$  and  $2 < p < 2_s^*$ . Later on, Shen and Yao [37] improved the corresponding results for the case  $\mu = 0$ . In the meanwhile, the authors in [41] disposed of the semiclassical ground state for the following fractional Schrödinger-Poisson system

$$\begin{cases} \varepsilon^{2s}(-\Delta)^s u + V(x)u + \phi u = f(u) + |u|^{2_s^*-2}u, & x \in \mathbb{R}^3, \\ \varepsilon^{2t}(-\Delta)^t \phi = u^2, & x \in \mathbb{R}^3. \end{cases}$$

Other meaningful results on fractional Schrödinger-Poisson system could be found in [14, 19, 22, 23, 32, 33, 36, 37, 39, 41, 44] and their references therein.

**1.2. Main results.** Motivated by all the works above, particularly by [19], we shall focus on the existence and concentration results for (1.1) with steep potential well. Because we are interested in positive solutions, without loss of generality, we assume that  $f \in C^0(\mathbb{R}, \mathbb{R})$  vanishes in  $(-\infty, 0)$  and satisfies the following conditions

- (f<sub>1</sub>)  $f \in C^0(\mathbb{R}, \mathbb{R}^+)$  and  $f(z) = o(z)$  as  $z \rightarrow 0$ , where  $\mathbb{R}^+ = [0, +\infty)$ ;
- (f<sub>2</sub>)  $|f(z)| \leq C_0(1 + |z|^{q-1})$  for some constants  $C_0 > 0$  and  $2 < q < 2_s^*$ ;
- (f<sub>3</sub>) there are some  $p \in \left(\frac{4s+2t}{s+t}, 2_s^*\right)$ ,  $\hat{\mu} > 0$  and  $\mu_0 > 0$  such that  $F(z) \geq \hat{\mu}z^p - \mu_0z^2$  for all  $z \in \mathbb{R}^+$ ;
- (f<sub>4</sub>) there is a  $\gamma > \frac{4s+2t}{s+t}$  such that  $zf(z) - \gamma F(z) \geq 0$  for all  $z \in \mathbb{R}^+$ , where  $F(z) = \int_0^z f(s)ds$ ;
- (f<sub>5</sub>) the map  $z \mapsto \frac{(s+t)f(z)z - 3F(z)}{z \frac{4s+2t}{s+t}}$  is nondecreasing on  $z \in (0, +\infty)$ .

Our first main result can be stated as follows.

**Theorem 1.1.** *Let  $s, t \in (0, 1)$  satisfy  $2s + 2t > 3$ . Suppose that  $(V_1)$ - $(V_3)$  and  $(f_1)$ - $(f_5)$  as well as the following conditions*

$(V_4)$   $V$  is weakly differentiable and  $(\nabla V, x) \in L^\infty(\mathbb{R}^3) \cup L^{\frac{3}{2s}}(\mathbb{R}^3)$  satisfies the inequality below

$$(s+t)(\gamma-2)V(x) + (x, \nabla V) \geq 0;$$

$(V_5)$  the map  $\theta \mapsto \theta^{2s}[(2s+2t-3)V(\theta x) - (\nabla V(\theta x), \theta x)]$  is nondecreasing on  $\theta \in (0, +\infty)$  and  $(2s+2t-3)V(x) \geq 2(\nabla V, x) \geq 0$  for all  $x \in \mathbb{R}^3$ .

If one of the following assumptions on  $p$  and  $\mu$  appearing in  $(f_3)$  holds true

$$(1.6) \quad \begin{cases} \text{(I)} : s > \frac{3}{4}, \frac{4s}{3-2s} < p < 2_s^* \text{ and for all } \hat{\mu} > 0; \\ \text{(II)} : s > \frac{3}{4}, \frac{4s+2t}{s+t} < p \leq \frac{4s}{3-2s} \text{ and for all sufficiently large } \hat{\mu} > 0; \\ \text{(III)} : \frac{1}{2} < s \leq \frac{3}{4}, \frac{4s+2t}{s+t} < p < 2_s^* \text{ and for all } \hat{\mu} > 0, \end{cases}$$

then there exists a  $\Lambda > 0$  such that the system (1.1) admits at least one positive ground state solution for all  $\lambda > \Lambda$ .

**Remark 1.2.** There exist many functions  $f$  that satisfy the assumptions  $(f_1)$ – $(f_5)$  above, for example  $f(z) = |z|^{\gamma-2}z$  for all  $z \in \mathbb{R}^+$  and  $f(z) = 0$  for all  $z < 0$ . Obviously, it would occur that  $\gamma < 4$  which results in some unpleasant difficulties. As to the potential  $V$ , without loss of generality, we are indeed assuming that it is of class  $\mathcal{C}^1$  at almost everywhere point in  $\mathbb{R}^3$  and provide an example as follows

$$V(x) = \begin{cases} 0, & \text{if } |x| \leq 1, \\ |x|^{\frac{2s+2t-3}{2}}, & \text{if } |x| > 1. \end{cases}$$

The reader is invited to infer that the restriction (1.6) is just used to restore the compactness. Moreover, we prefer to believe that the example on  $V$  above is not sharp, but it reveals that the existence result in Theorem 1.1 seems reasonable.

Inspired by the results in [5, 45], we get the following concentration result:

**Theorem 1.3.** *Let  $(u_\lambda, \phi_{u_\lambda}) \in H^s(\mathbb{R}^3) \times D^{t,2}(\mathbb{R}^3)$  be the ground state solution obtained by Theorem 1.1, then  $u_\lambda \rightarrow u_0$  in  $H^s(\mathbb{R}^3)$  and  $\phi_{u_\lambda} \rightarrow \phi_{u_0}$  in  $D^{t,2}(\mathbb{R}^3)$  along a subsequence as  $\lambda \rightarrow +\infty$ , where  $u_0 \in H_0^s(\Omega)$  is a ground state solution of*

$$(1.7) \quad \begin{cases} (-\Delta)^s u + c_t \left( \int_\Omega \frac{u^2(y)}{|x-y|^{3-2t}} dy \right) u = f(u) + |u|^{2_s^*-2} u, & \text{in } \Omega, \\ u = 0, & \text{on } \partial\Omega. \end{cases}$$

Here  $c_t > 0$  is a constant given by (2.8) below.

As pointed out in Remark 1.2, the assumptions on  $f$  and  $V$  required in Theorem 1.1 are somehow restrictive. It is natural to ask that whether the existence result remains true when  $(f_5)$  and  $(V_5)$  are absent. Thus, our next main result shows an affirmative answer.

**Theorem 1.4.** *Let  $s, t \in (0, 1)$  satisfy  $2s + 2t > 3$ . Suppose that  $(V_1)$ - $(V_4)$  and  $(f_1)$ - $(f_4)$ . If one of the assumptions in (1.6) holds true, then there exists a  $\hat{\Lambda} > 0$  such that the system (1.1) has at least one positive solution for all  $\lambda > \hat{\Lambda}$ .*

**Remark 1.5.** It is worth pointing out here that even if we only consider the case  $s = t = 1$  in Theorem 1.4, in contrast to [45, Theorem 1.3], there are three main contributions:

- (1) Firstly, the more general nonlinearity is dealt with and it needs some more careful calculations;
- (2) Secondly, the critical term in the nonlinearity is involved and so we have to take some deep and delicate analysis to restore the compactness;

(3) Last but not the least, we do not assume a weight function  $K$  in the front of the Poisson term in (1.1). Actually, if we follow the arguments adopted in this quoted paper, the weight function  $K$  with  $K \in L^{\frac{6}{4s+2t-3}}(\mathbb{R}^3)$  seems indispensable. So, we can relax the constraint assumption in this direction.

Proceeding as the same way in Theorem 1.3, we can also derive the asymptotical behavior of solutions obtained in Theorem 1.4. More precisely, we shall demonstrate the following result whose detailed proof is omitted.

**Theorem 1.6.** *Let  $(u_\lambda, \phi_{u_\lambda}) \in H^s(\mathbb{R}^3) \times D^{t,2}(\mathbb{R}^3)$  denote the positive solution in Theorem 1.4, then  $u_\lambda \rightarrow u_0$  in  $H^s(\mathbb{R}^3)$  and  $\phi_{u_\lambda} \rightarrow \phi_{u_0}$  in  $D^{t,2}(\mathbb{R}^3)$  along a subsequence as  $\lambda \rightarrow +\infty$ , where  $u_0 \in H_0^s(\Omega)$  is a positive solution of (1.7)*

As far as we are concerned, the main results in this article seem new by now. Alternatively, it should be mentioned that this paper could be regarded as a continuation of our latest work in [35], where the existence and concentrating results of planar Schrödinger-Poisson equation with steep potential well were established. Whereas, there are two essential differences: On the one hand, due to the different geometry structures of the two variational functionals, we must take advantage of some new techniques to restore the compactness; On the other hand, since we consider the existence of ground state solutions in Theorem 1.1, a suitable constraint minimization argument will be used, instead of depending on the mountain-pass theorem in [35]. Finally, when the critical term  $|u|^{2^*_s-2}u$  in the systme (1.1) disappears, one may be curious about the case that the potential is strongly indefinite according to [26]. Of course, we are also working hard in this direction and it would be contemplated in our further studies.

The paper is organized as follows. In Section 2, we mainly introduce some preliminary results. In Sections 3 and 4, we show some crucial lemmas and exhibit the detailed proofs of Theorems 1.1, 1.3 and 1.4, respectively.

**Notations:** From now on in this paper, otherwise mentioned, we utilize the following notations:

- $C, C_1, C_2, \dots$  denote any positive constant, whose value is not relevant and  $\mathbb{R}^+ \triangleq (0, +\infty)$ .
- Let  $(Z, \|\cdot\|_Z)$  be a Banach space with dual space  $(Z^{-1}, \|\cdot\|_{Z^{-1}})$ , and  $\Phi$  be functional on  $Z$ .
- The (PS) sequence at a level  $c \in \mathbb{R}$  ((PS) $_c$  sequence in short) corresponding to  $\Phi$  means that  $\Phi(x_n) \rightarrow c$  and  $\Phi'(x_n) \rightarrow 0$  in  $Z^{-1}$  as  $n \rightarrow \infty$ , where  $\{x_n\} \subset Z$ .
- $\|\cdot\|_p$  stands for the usual norm of the Lebesgue space  $L^p(\mathbb{R}^N)$  for all  $p \in [1, +\infty]$ , and  $\|\cdot\|_{H^\alpha(\mathbb{R}^N)}$  denotes the usual norm of the Sobolev space  $H^\alpha(\mathbb{R}^N)$  for  $\alpha \in (0, 1)$ .
- For any  $\varrho > 0$  and every  $x \in \mathbb{R}^3$ ,  $B_\varrho(x) \triangleq \{y \in \mathbb{R}^3 : |y - x| < \varrho\}$ .
- $o_n(1)$  denotes the real sequences with  $o_n(1) \rightarrow 0$  as  $n \rightarrow +\infty$ .
- “ $\rightarrow$ ” and “ $\rightharpoonup$ ” stand for the strong and weak convergence in the related function spaces, respectively;

## 2. PRELIMINARY STUFF

**2.1. Variational setting.** In this section, according to the explorations about the fractional Sobolev spaces in [25], we first bring in some necessary variational settings which permit us to treat the problems variationally. Denote the fractional Sobolev space  $W^{\alpha,p}(\mathbb{R}^N)$  for any  $p \in [1, +\infty)$  and  $\alpha \in (0, 1)$  by

$$W^{\alpha,p}(\mathbb{R}^N) = \left\{ u \in L^p(\mathbb{R}^N) : \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^p}{|x - y|^{N+\alpha p}} dx dy < +\infty \right\}$$

equipped with the natural norm

$$\|u\|_{W^{\alpha,p}(\mathbb{R}^N)} = \left( \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^p}{|x - y|^{N+\alpha p}} dx dy + \int_{\mathbb{R}^N} |u|^p dx \right)^{\frac{1}{p}}.$$

In particular, the fractional Sobolev space  $W^{\alpha,2}(\mathbb{R}^N)$  would be simply relabeled by  $H^\alpha(\mathbb{R}^N)$  if  $p = 2$ . As a matter of fact, the Hilbert space  $H^\alpha(\mathbb{R}^N)$  can also be described by the Fourier transform, that is,

$$H^\alpha(\mathbb{R}^N) = \left\{ u \in L^2(\mathbb{R}^N) : \int_{\mathbb{R}^N} |\xi|^{2\alpha} |\widehat{u}(\xi)|^2 + |\widehat{u}(\xi)|^2 d\xi < +\infty \right\},$$

where  $\widehat{u}$  denotes the usual Fourier transform of  $u$ . When we take the definition of the fractional Sobolev space  $H^\alpha(\mathbb{R}^N)$  by the Fourier transform, the inner product and the norm for  $H^\alpha(\mathbb{R}^N)$  are defined as

$$(u, v)_{H^\alpha(\mathbb{R}^N)} = \int_{\mathbb{R}^N} |\xi|^{2\alpha} \widehat{u}(\xi) \widehat{v}(\xi) + \widehat{u}(\xi) \widehat{v}(\xi) d\xi, \quad \forall u, v \in H^\alpha(\mathbb{R}^N).$$

and

$$\|u\|_{H^\alpha(\mathbb{R}^N)} = \left( \int_{\mathbb{R}^N} |\xi|^{2\alpha} |\widehat{u}(\xi)|^2 + |\widehat{u}(\xi)|^2 d\xi \right)^{\frac{1}{2}}, \quad \forall u, v \in H^\alpha(\mathbb{R}^N).$$

Thanks to the Plancherel's theorem, we have  $\|u\|_2 = \|\widehat{u}\|_2$  and  $\|(-\Delta)^{\frac{\alpha}{2}} u\|_2 = \|\xi|^\alpha \widehat{u}\|_2$ . Hence

$$(2.1) \quad \|u\|_{H^\alpha(\mathbb{R}^N)} = \left( \int_{\mathbb{R}^N} |(-\Delta)^{\frac{\alpha}{2}} u|^2 + |u|^2 dx \right)^{\frac{1}{2}}, \quad \forall u \in H^\alpha(\mathbb{R}^N).$$

It infers from [25, Proposition 3.4 and Proposition 3.6] that

$$\|(-\Delta)^{\frac{\alpha}{2}} u\|_2 = \left( \int_{\mathbb{R}^N} |\xi|^{2\alpha} |\widehat{u}(\xi)|^2 d\xi \right)^{\frac{1}{2}} = \left( \frac{1}{C_N(\alpha)} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^2}{|x - y|^{N+2\alpha}} dx dy \right)^{\frac{1}{2}}.$$

showing that the norm in (2.1) makes sense for the fractional Sobolev space. Moreover, we introduce the homogeneous fractional Sobolev space  $D^{\alpha,2}(\mathbb{R}^N)$  by

$$D^{\alpha,2}(\mathbb{R}^N) = \left\{ u \in L^{2^*_\alpha}(\mathbb{R}^N) : |\xi|^\alpha \widehat{u}(\xi) \in L^{2^*_\alpha}(\mathbb{R}^N) \right\} \text{ with } 2^*_\alpha = \frac{2N}{N-2\alpha} \text{ and } N \geq 3,$$

which is the completion of  $C_0^\infty(\mathbb{R}^N)$  under the norm

$$\|u\|_{D^{\alpha,2}(\mathbb{R}^N)} = \left( \int_{\mathbb{R}^N} |(-\Delta)^{\frac{\alpha}{2}} u|^2 dx \right)^{\frac{1}{2}} = \left( \int_{\mathbb{R}^N} |\xi|^{2\alpha} |\widehat{u}(\xi)|^2 d\xi \right)^{\frac{1}{2}}, \quad \forall u \in D^{\alpha,2}(\mathbb{R}^N).$$

Taking into account the imbedding theorem  $H^\alpha(\mathbb{R}^N) \hookrightarrow L^r(\mathbb{R}^N)$  for every  $r \in [2, 2^*_\alpha)$ , there exists a constant  $C_r > 0$  such that

$$(2.2) \quad \|u\|_{H^\alpha(\mathbb{R}^N)} \leq C_r \|u\|_r, \quad \forall u \in H^\alpha(\mathbb{R}^N) \text{ and } 2 \leq r < 2^*_\alpha.$$

Also there exists a best constant  $S_\alpha > 0$  (see e.g. [13]) such that

$$(2.3) \quad S_\alpha = \inf_{u \in D^{\alpha,2}(\mathbb{R}^N) \setminus \{0\}} \frac{\int_{\mathbb{R}^N} |(-\Delta)^{\frac{\alpha}{2}} u|^2 dx}{\left( \int_{\mathbb{R}^N} |u|^{2^*_\alpha} dx \right)^{\frac{2}{2^*_\alpha}}}.$$

Throughout this paper, for  $s \in (0, 1)$  and the dimension  $N = 3$ , we define the space

$$E \triangleq \left\{ u \in H^s(\mathbb{R}^3) : \int_{\mathbb{R}^3} V(x) u^2 dx < +\infty \right\}.$$

By using  $(V_1)$ , it is easy to verify that it is a Hilbert space equipped with the inner product and norm

$$(u, v)_E = \int_{\mathbb{R}^3} (-\Delta)^{\frac{s}{2}} u (-\Delta)^{\frac{s}{2}} v + V(x) uv dx \text{ and } \|u\|_E = \left( \int_{\mathbb{R}^3} |(-\Delta)^{\frac{s}{2}} u|^2 + V(x) u^2 dx \right)^{\frac{1}{2}}$$

for any  $u, v \in E$ . Particularly, one can deduce that the imbedding  $E \hookrightarrow H^s(\mathbb{R}^3)$  is continuous. Indeed, combining (V<sub>2</sub>) and (2.3), one has

$$\begin{aligned} \int_{\mathbb{R}^3} u^2 dx &= \int_{\mathbb{R}^3 \setminus \Sigma} u^2 dx + \int_{\Sigma} u^2 dx \leq \frac{1}{c} \int_{\mathbb{R}^3 \setminus \Sigma} V(x) u^2 dx + |\Sigma|^{\frac{2_s^* - 2}{2_s^*}} \left( \int_{\Sigma} |u|^{2_s^*} dx \right)^{\frac{2}{2_s^*}} \\ &\leq \max \left\{ \frac{1}{c}, |\Sigma|^{\frac{2_s^* - 2}{2_s^*}} \right\} \|u\|_E^2, \end{aligned}$$

where  $|\Sigma|$  stands for the Lebesgue measure of a Lebesgue measurable set  $\Sigma \subset \mathbb{R}^3$ . As a consequence of (2.2) and (2.3), there exists a constant  $d_r > 0$  such that

$$(2.4) \quad |u|_r \leq d_r \|u\|_E, \quad \forall u \in E \text{ and } 2 \leq r \leq 2_s^*.$$

For any  $\lambda > 0$ , define the Hilbert space  $E_\lambda \triangleq (E, \|\cdot\|_{E_\lambda})$  with inner product and norm given by

$$(u, v)_{E_\lambda} = \int_{\mathbb{R}^3} (-\Delta)^{\frac{s}{2}} u (-\Delta)^{\frac{s}{2}} v + \lambda V(x) u v dx \text{ and } \|u\|_{E_\lambda} = \left( \int_{\mathbb{R}^3} |(-\Delta)^{\frac{s}{2}} u|^2 + \lambda V(x) |u|^2 dx \right)^{\frac{1}{2}}$$

for all  $u, v \in E$ . Obviously, if  $\lambda \geq 1$ , one sees  $\|u\|_E \leq \|u\|_{E_\lambda}$  for all  $u \in E$ . Using (V<sub>2</sub>) again,

$$\begin{cases} \int_{\Sigma} |u|^2 dx \leq |\Sigma|^{\frac{2_s^* - 2}{2_s^*}} |u|_{2_s^*}^2 \leq |\Sigma|^{\frac{2_s^* - 2}{2_s^*}} S_s^{-1} \|u\|_{E_\lambda}^2, \\ \int_{\mathbb{R}^3 \setminus \Sigma} |u|^2 dx \leq \frac{1}{\lambda c} \int_{\mathbb{R}^3 \setminus \Sigma} \lambda V(x) |u|^2 dx \leq \frac{1}{\lambda c} \int_{\mathbb{R}^3} \lambda V(x) |u|^2 dx \leq \frac{1}{\lambda c} \|u\|_{E_\lambda}^2. \end{cases}$$

From which, for any  $r \in [2, 2_s^*]$ , there holds

$$\begin{aligned} \int_{\mathbb{R}^3} |u|^r dx &\leq \left( \int_{\mathbb{R}^3} |u|^2 dx \right)^{\frac{2_s^* - r}{2_s^* - 2}} \left( \int_{\mathbb{R}^3} |u|^{2_s^*} dx \right)^{\frac{r - 2}{2_s^* - 2}} \\ &\leq \left( \max \left\{ S_s^{-1} |\Sigma|^{\frac{2_s^* - 2}{2_s^*}}, \frac{1}{\lambda c} \right\} \|u\|_{E_\lambda}^2 \right)^{\frac{2_s^* - r}{2_s^* - 2}} \left( S_s^{-\frac{2_s^*}{2}} \|u\|_{E_\lambda}^{2_s^*} \right)^{\frac{r - 2}{2_s^* - 2}}. \end{aligned}$$

Hence, for all  $r \in [2, 2_s^*]$ , we reach

$$(2.5) \quad \int_{\mathbb{R}^3} |u|^r dx \leq |\Sigma|^{\frac{2_s^* - r}{2_s^*}} S_s^{-\frac{r}{2}} \|u\|_{E_\lambda}^r \text{ whenever } \lambda \geq c^{-1} |\Sigma|^{-\frac{2_s^* - 2}{2_s^*}} S_s.$$

When the work space  $E_\lambda$  is built, we turn to find the variational structure of system (1.1). Following the classic Schrödinger-Poisson system, it can reduce to be a single equation. Actually, according to the Hölder's inequality, for every  $u \in H^s(\mathbb{R}^3)$  and  $v \in D^{t,2}(\mathbb{R}^3)$ , one has

$$(2.6) \quad \begin{aligned} \int_{\mathbb{R}^3} u^2 v dx &\leq \left( \int_{\mathbb{R}^3} |u|^{\frac{12}{3+2t}} dx \right)^{\frac{3+2t}{6}} \left( \int_{\mathbb{R}^3} |v|^{\frac{6}{3-2t}} dx \right)^{\frac{3-2t}{6}} \\ &\leq S_t^{-\frac{1}{2}} \|u\|_{H^s(\mathbb{R}^3)}^2 \|v\|_{D^{t,2}(\mathbb{R}^3)} \leq C \|u\|_{H^s(\mathbb{R}^3)}^2 \|v\|_{D^{t,2}(\mathbb{R}^3)}, \end{aligned}$$

where we have used the continuous imbedding  $H^s(\mathbb{R}^3) \hookrightarrow L^{\frac{12}{3+2t}}(\mathbb{R}^3)$  since  $4s + 2t > 3$  and  $t \in (0, 1)$ .

Given  $u \in H^s(\mathbb{R}^3)$ , one can use the Lax-Milgram theorem and then there exists a unique  $\phi_u^t \in D^{t,2}(\mathbb{R}^3)$  such that

$$(2.7) \quad \int_{\mathbb{R}^3} (-\Delta)^t \phi_u^t v dx = \int_{\mathbb{R}^3} (-\Delta)^{\frac{t}{2}} \phi_u^t (-\Delta)^{\frac{t}{2}} v dx = \int_{\mathbb{R}^3} u^2 v dx, \quad \forall v \in D^{t,2}(\mathbb{R}^3),$$

showing that  $\phi_u^t$  satisfies the Poisson equation

$$(-\Delta)^t \phi_u^t = u^2, \quad x \in \mathbb{R}^3.$$

In view of [25], its integral expression can be characterized by

$$(2.8) \quad \phi_u^t(x) = c_t \int_{\mathbb{R}^3} \frac{u^2(y)}{|x-y|^{3-2t}} dx, \quad x \in \mathbb{R}^3,$$

which is called  $t$ -Riesz potential, where

$$c_t = \pi^{-\frac{3}{2}} 2^{-2t} \frac{\Gamma(\frac{3}{2} - 2t)}{\Gamma(t)}.$$

It follows from (2.8) that  $\phi_u^t(x) \geq 0$  for all  $x \in \mathbb{R}^3$ . Taking  $v = \phi_u^t$  in (2.6) and (2.7), we derive

$$(2.9) \quad \|\phi_u^t\|_{D^{t,2}(\mathbb{R}^3)} \leq C \|u\|_{H^s(\mathbb{R}^3)}^2.$$

Substituting (2.8) into (1.1), one can rewrite (1.1) in the following equivalent form

$$(2.10) \quad (-\Delta)^s u + \lambda V(x)u + \phi_u^t u = f(u) + |u|^{2_s^*-2} u, \quad x \in \mathbb{R}^3.$$

The variational functional  $I_\lambda : E_\lambda \rightarrow \mathbb{R}$  associated to the problem (2.10) is given by

$$(2.11) \quad I_\lambda(u) = \frac{1}{2} \|u\|_{E_\lambda}^2 + \frac{1}{4} \int_{\mathbb{R}^3} \phi_u^t u^2 dx - \int_{\mathbb{R}^3} F(u) dx - \frac{1}{2_s^*} \int_{\mathbb{R}^3} |u|^{2_s^*} dx.$$

It would be simply verified that  $I_\lambda$  is well-defined in  $E_\lambda$  and belongs to  $\mathcal{C}^1(E_\lambda, \mathbb{R})$  whose derivative is given by

$$I'_\lambda(u)v = \int_{\mathbb{R}^3} [(-\Delta)^{\frac{s}{2}} u (-\Delta)^{\frac{s}{2}} v + \lambda V(x)uv] dx + \int_{\mathbb{R}^3} \phi_u^t uv dx - \int_{\mathbb{R}^3} (f(u) + |u|^{2_s^*-2} u) v dx$$

for any  $u, v \in E_\lambda$ . It is clear to see that if  $u$  is a critical point of  $I_\lambda$ , then the pair  $(u, \phi_u^t)$  is a solution of system (1.1).

**2.2. Basic lemmas.** It is similar to the proof of [39, Proposition 2.1] that we can derive the following

**Lemma 2.1.** (Pohožaev identity) *Let  $u \in E_\lambda$  be a critical point of the functional  $I_\lambda$ , then the identity  $P_\lambda(u) \equiv 0$  holds true, where the functional  $P_\lambda : E_\lambda \rightarrow \mathbb{R}$  is defined by*

$$P_\lambda(u) \triangleq \frac{3-2s}{2} \int_{\mathbb{R}^3} |(-\Delta)^{\frac{s}{2}} u|^2 dx + \frac{1}{2} \int_{\mathbb{R}^3} [3V(x) + (\nabla V, x)] |u|^2 dx + \frac{2t+3}{4} \int_{\mathbb{R}^3} \phi_u^t u^2 dx \\ - 3 \int_{\mathbb{R}^3} F(u) dx - \frac{3}{2_s^*} \int_{\mathbb{R}^3} |u|^{2_s^*} dx.$$

Now, let us define the functional  $N : E_\lambda \rightarrow \mathbb{R}$  by

$$N(u) = \int_{\mathbb{R}^3} \phi_u^t u^2 dx, \quad \forall u \in E_\lambda.$$

We gather the results in [37, Lemmas 9 and 10] to introduce the properties associated with  $N$  below.

**Lemma 2.2.** *Let  $s, t \in (0, 1)$  satisfy  $4s + 2t > 3$ , then the following properties are true:*

- (1) *For all  $u \in E_\lambda$  and we set  $u_\theta(\cdot) \triangleq \theta^{s+t} u(\theta \cdot)$  for  $\theta \in \mathbb{R}^+$ , then  $N(u_\theta) = \theta^{4s+2t-3} N(u)$ ;*
- (2)  *$\phi_{u(\cdot+y)}^t = \phi_u^t(\cdot + y)$  for all  $y \in \mathbb{R}^3$ .*
- (3) *If  $u_n \rightharpoonup u$  in  $E_\lambda$ , then  $N(u_n) - N(u_n - u) - N(u) = o_n(1)$  in  $E_\lambda$ ,  $N'(u_n) - N'(u_n - u) - N'(u) = o_n(1)$  in  $(E_\lambda)^{-1}$ .*

We conclude this section by the following Vanishing lemma associated with the fractional Sobolev space.



**Lemma 2.3.** (See e.g. [31, Lemma]) Assume  $(u_n)$  is a bounded sequence in  $H^\alpha(\mathbb{R}^3)$  with  $\alpha \in (0, 1)$ . If

$$\lim_{n \rightarrow \infty} \sup_{y \in \mathbb{R}^3} \int_{B_\varrho(y)} |u_n|^2 dx = 0$$

for some  $\varrho > 0$ , then  $u_n \rightarrow 0$  in  $L^q(\mathbb{R}^3)$  for all  $2 < p < 2_\alpha^*$ .

### 3. EXISTENCE AND CONCENTRATION

In this section, we focus on the existence and concentration of ground state solutions for (1.1). First of all, to look for a ground state solution, we shall consider the following minimization problem

$$(3.1) \quad m_\lambda \triangleq \inf_{u \in \mathcal{M}_\lambda} I_\lambda(u),$$

where  $\mathcal{M}_\lambda = \{u \in E_\lambda \setminus \{0\} : G_\lambda(u) = 0\}$  with the functional  $G_\lambda : E_\lambda \rightarrow \mathbb{R}$  defined by

$$\begin{aligned} G_\lambda(u) &= \frac{4s+2t-3}{2} \int_{\mathbb{R}^3} |(-\Delta)^{\frac{s}{2}} u|^2 dx + \frac{1}{2} \int_{\mathbb{R}^3} \lambda[(2s+2t-3)V(x) - (\nabla V, x)]u^2 dx \\ &\quad + \frac{4s+2t-3}{4} \int_{\mathbb{R}^3} \phi_u^t u^2 dx - \int_{\mathbb{R}^3} [(s+t)f(u)u - 3F(u)]dx - \frac{2_s^*(s+t)-3}{2_s^*} \int_{\mathbb{R}^3} |u|^{2_s^*} dx. \end{aligned}$$

Recalling the functional  $P_\lambda$  in Lemma 2.1, one sees that  $G_\lambda(u) = (s+t)I'_\lambda(u)u - P_\lambda(u)$  for all  $u \in E_\lambda$ . In other words, if  $u \in E_\lambda$  is a critical point of  $I_\lambda$ , then we are derived from Lemma 2.1 that  $G_\lambda(u) = 0$ . As a consequence, the set  $\mathcal{M}_\lambda$  is a natural constraint and we then begin showing some properties for it and the minimization constant  $m_\lambda$ .

Before exhibiting them, we need the following elementary facts:

$$(3.2) \quad \xi(\theta, x) \triangleq V(x) - \theta^{2s+2t-3}V(\theta^{-1}x) - \frac{1 - \theta^{4s+2t-3}}{4s+2t-3} [(2s+2t-3)V(x) - (\nabla V, x)] \geq 0$$

for all  $(\theta, x) \in (0, +\infty) \times \mathbb{R}^3$  and

$$(3.3) \quad \zeta(\theta, z) \triangleq \frac{1 - \theta^{4s+2t-3}}{4s+2t-3} [(s+t)f(z)z - 3F(z)] + \theta^{-3}F(\theta^{s+t}z) - F(z) \geq 0$$

for all  $(\theta, z) \in (0, +\infty) \times \mathbb{R}^+$ .

Actually, since  $V$  is weakly differentiable by  $(V_4)$ , one uses  $(V_5)$  to see that

$$\begin{aligned} \frac{\partial}{\partial \theta} \xi(\theta, x) &= \theta^{4s+2t-4} \left\{ [(2s+2t-3)V(x) - (\nabla V, x)] - \frac{(2s+2t-3)V(\theta^{-1}x) - (\nabla V(\theta^{-1}x), \theta^{-1}x)}{\theta^{2s}} \right\} \\ &\begin{cases} \leq 0, & \text{if } \theta \in (0, 1] \\ \geq 0, & \text{if } \theta \in [1, +\infty). \end{cases} \end{aligned}$$

Hence, the function  $\theta \mapsto \xi(\theta, x)$  is decreasing on  $(0, 1)$  and increasing on  $(1, +\infty)$  for all  $x \in \mathbb{R}^3$  which indicate that  $\xi(\theta, x) \geq \min_{\theta > 0} \xi(\theta, x) = \xi(1, x) = 0$  for all  $(\theta, x) \times (0, +\infty) \in \mathbb{R}^3$ . Similarly, we are able to

apply  $(f_5)$  to derive

$$\begin{aligned} \frac{\partial}{\partial \theta} \zeta(\theta, z) &= \theta^{-4} [(s+t)f(\theta^{s+t}z)\theta^{s+t}z - 3F(\theta^{s+t}z)] - \theta^{4s+2t-4} [(s+t)f(z)z - 3F(z)] \\ &= \theta^{4s+2t-4} z^{\frac{4s+2t}{s+t}} \left[ \frac{(s+t)f(\theta^{s+t}z)\theta^{s+t}z - 3F(\theta^{s+t}z)}{(\theta^{s+t}z)^{\frac{4s+2t}{s+t}}} - \frac{(s+t)f(z)z - 3F(z)}{z^{\frac{4s+2t}{s+t}}} \right] \\ &\begin{cases} \leq 0, & \text{if } \theta \in (0, 1] \\ \geq 0, & \text{if } \theta \in [1, +\infty). \end{cases} \end{aligned}$$

It therefore infers that  $\zeta(\theta, z) \geq \min_{\theta > 0} \zeta(\theta, z) = \zeta(1, z) = 0$  for all  $(\theta, z) \times (0, +\infty) \in \mathbb{R}^+$ .

**Lemma 3.1.** *Let  $s, t \in (0, 1)$  satisfy  $2s + 2t > 3$ . Assume  $(V_1) - (V_3)$  with  $(V_4) - (V_5)$  and  $(f_1) - (f_3)$  with  $(f_5)$ , then for any nonzero  $u \in E_\lambda$ , there is a unique  $\bar{\theta} = \bar{\theta}(u) > 0$  such that  $u_{\bar{\theta}} = \bar{\theta}^{s+t}u(\bar{\theta} \cdot) \in \mathcal{M}_\lambda$  for suitably large  $\lambda > 0$ , where  $I_\lambda(u_{\bar{\theta}}) = \max_{\theta > 0} I_\lambda(u_\theta)$ . In particular, there holds*

$$m_\lambda = d_\lambda \triangleq \inf_{u \in E_\lambda \setminus \{0\}} \max_{\theta > 0} I_\lambda(u_\theta).$$

*Proof.* For any  $u \in E_\lambda \setminus \{0\}$  and  $\theta > 0$ , we define  $\tau(\theta) = I_\lambda(u_\theta)$ , where

$$\begin{aligned} \tau(\theta) &= \frac{\theta^{4s+2t-3}}{2} \int_{\mathbb{R}^3} |(-\Delta)^{\frac{s}{2}} u|^2 dx + \frac{\theta^{2s+2t-3}}{2} \int_{\mathbb{R}^3} \lambda V(\theta^{-1}x) u^2 dx + \frac{\theta^{4s+2t-3}}{4} \int_{\mathbb{R}^3} \phi_u^t u^2 dx \\ &\quad - \theta^{-3} \int_{\mathbb{R}^3} F(\theta^{s+t}u) dx - \frac{\theta^{2_s^*(s+t)-3}}{2_s^*} \int_{\mathbb{R}^3} |u|^{2_s^*} dx. \end{aligned}$$

It is simple to observe that

$$\tau'(\theta) = 0 \iff \theta^{-1}G_\lambda(u_\theta) = 0 \iff G_\lambda(u_\theta) = 0 \iff u_\theta \in \mathcal{M}_\lambda.$$

Since  $4s + 2t < 2_s^*(s + t)$  and  $\lim_{\theta \rightarrow 0^+} \theta^{-3}F(\theta^{s+t}z) = 0$  for all  $z \in \mathbb{R}$  by  $(f_3)$ , we can derive  $\lim_{\theta \rightarrow 0^+} \tau(\theta) > 0$ . Without loss of generality, we are assuming that  $0 \in \Omega$  in  $(V_3)$  and thus  $\lim_{\theta \rightarrow +\infty} \int_{\mathbb{R}^3} \lambda V(\theta^{-1}x) u^2 dx = 0$ . Adopting  $4s + 2t < 2_s^*(s + t)$  and  $(f_3)$  again, it holds that  $\lim_{\theta \rightarrow +\infty} \tau(\theta) = -\infty$ . As a consequence, with the above two facts in hands, we take advantage of  $4s + 2t < 2_s^*(s + t)$  and  $(f_3)$  to demonstrate that  $\tau(\theta)$  possesses a critical point which corresponds to its maximum, that is, there exists a constant  $\bar{\theta} > 0$  such that  $\tau'(\bar{\theta}) = 0$ . We next verify that  $\bar{\theta}$  is unique. Arguing it indirectly, we would assume that there exist two constants  $\theta_1, \theta_2 > 0$  with  $\theta_1 \neq \theta_2$  such that  $u_{\theta_i} \in \mathcal{M}_\lambda$  for  $i \in \{1, 2\}$ . It concludes from some elementary computations that

$$\begin{aligned} I_\lambda(u_{\theta_1}) - I_\lambda(u_{\theta_2}) &- \frac{\theta_1^{4s+2t-3} - \theta_2^{4s+2t-3}}{(4s + 2t - 3)\theta_1^{4s+2t-3}} G_\lambda(u_{\theta_1}) \\ &= \frac{\theta_1^{2s+2t-3}}{2} \int_{\mathbb{R}^3} \xi \left( \frac{\theta_2}{\theta_1}, \theta_1 x \right) u^2 dx + \theta_1^{-3} \int_{\mathbb{R}^3} \zeta \left( \frac{\theta_2}{\theta_1}, \theta_1^{s+t}u \right) dx \\ &\quad + \frac{\theta_1^{2_s^*(s+t)-3}}{2_s^*} \left[ \frac{1 - \left( \frac{\theta_2}{\theta_1} \right)^{4s+2t-3}}{4s + 2t - 3} [2_s^*(s + t) - 3] + \left( \frac{\theta_2}{\theta_1} \right)^{2_s^*(s+t)-3} - 1 \right] \int_{\mathbb{R}^3} |u|^{2_s^*} dx \end{aligned}$$

and

$$\begin{aligned} I_\lambda(u_{\theta_2}) - I_\lambda(u_{\theta_1}) &- \frac{\theta_2^{4s+2t-3} - \theta_1^{4s+2t-3}}{(4s + 2t - 3)\theta_2^{4s+2t-3}} G_\lambda(u_{\theta_2}) \\ &= \frac{\theta_2^{2s+2t-3}}{2} \int_{\mathbb{R}^3} \xi \left( \frac{\theta_1}{\theta_2}, \theta_2 x \right) u^2 dx + \theta_2^{-3} \int_{\mathbb{R}^3} \zeta \left( \frac{\theta_1}{\theta_2}, \theta_2^{s+t}u \right) dx \\ &\quad + \frac{\theta_2^{2_s^*(s+t)-3}}{2_s^*} \left[ \frac{1 - \left( \frac{\theta_1}{\theta_2} \right)^{4s+2t-3}}{4s + 2t - 3} [2_s^*(s + t) - 3] + \left( \frac{\theta_1}{\theta_2} \right)^{2_s^*(s+t)-3} - 1 \right] \int_{\mathbb{R}^3} |u|^{2_s^*} dx. \end{aligned}$$

In view of (3.2) and (3.3), combining the above two formulas with  $G_\lambda(u_{\theta_i}) = 0$  for  $i \in \{1, 2\}$ , we arrive at a contradiction if  $\theta_1 \neq \theta_2$ . Finally, the result  $d_\lambda \leq m_\lambda$  is a direct consequence of the inequality

$$(3.4) \quad I_\lambda(u) - I_\lambda(u_\theta) - \frac{1 - \theta^{4s+2t-3}}{4s + 2t - 3} G_\lambda(u) \geq 0, \quad \forall u \in E_\lambda \text{ and } \theta > 0,$$

we immediately finish the proof of this lemma.  $\square$

The following results can be found in [40].

**Lemma 3.2.** *Let  $u_\varepsilon$  be defined by (3.10) in the proof of Lemma 3.3 below, then*

$$(3.5) \quad \int_{\mathbb{R}^3} |(-\Delta)^{\frac{s}{2}} u_\varepsilon|^2 dx \leq S_s^{\frac{3}{2s}} + O(\varepsilon^{3-2s}),$$

and

$$(3.6) \quad \int_{\mathbb{R}^3} |u_\varepsilon|^{2^*} dx = S_s^{\frac{3}{2s}} + O(\varepsilon^3).$$

For all  $q \in [2, 2_s^*)$ , there holds

$$(3.7) \quad \int_{\mathbb{R}^3} |u_\varepsilon|^q dx = \begin{cases} O\left(\varepsilon^{3-\frac{3-2s}{2}q}\right), & \text{for } q > \frac{3}{3-2s}, \\ O\left(\varepsilon^{\frac{3}{2}} |\log \varepsilon|\right), & \text{for } q = \frac{3}{3-2s}, \\ O\left(\varepsilon^{\frac{3-2s}{2}q}\right), & \text{for } q < \frac{3}{3-2s}. \end{cases}$$

According to Lemma 3.1, we know that  $\mathcal{M}_\lambda$  is a nonempty set for some suitably large  $\lambda > 0$ . The following lemma ensures that the minimization constant  $m_\lambda$  would be well-defined. More precisely, we further show that  $m_\lambda$  is uniformly bounded from below and above by some positive constants which are independent of some suitably large  $\lambda > 0$ .

**Lemma 3.3.** *Let  $s, t \in (0, 1)$  satisfy  $2s + 2t > 3$ . Assume that  $(V_1) - (V_5)$  and  $(f_1) - (f_5)$ , there is a  $\rho > 0$  independent of  $\lambda > \Lambda_0$  such that*

$$(3.8) \quad \inf_{\lambda > \Lambda_0} m_\lambda \geq \rho,$$

where  $\Lambda_0 \triangleq \max\{1, c^{-1}|\Sigma|^{-\frac{2_s^*-2}{2_s^*}} S_s\}$ . If in addition one of the assumptions in (1.6) holds true, then

$$(3.9) \quad \sup_{\lambda > \Lambda_0} m_\lambda < \frac{s}{3} S_s^{\frac{3}{2s}}.$$

*Proof.* For all  $u \in \mathcal{M}_\lambda$ , we are derived from  $(f_4)$  and  $(\nabla V, x) \geq 0$  for all  $x \in \mathbb{R}^3$  in  $(V_5)$  that

$$\begin{aligned} I_\lambda(u) &= I_\lambda(u) - \frac{1}{(s+t)\gamma-3} G_\lambda(u) \\ &= \frac{(s+t)\gamma - (4s+2t)}{2[(s+t)\gamma-3]} |(-\Delta)^{\frac{s}{2}} u|_2^2 + \frac{1}{2[(s+t)\gamma-3]} \int_{\mathbb{R}^3} \lambda [(s+t)(\gamma-2)V(x) + (\nabla V, x)] u^2 dx \\ &\quad + \frac{(s+t)\gamma - (4s+2t)}{4[(s+t)\gamma-3]} \int_{\mathbb{R}^3} \phi_u^t u^2 dx + \frac{s+t}{(s+t)\gamma-3} \int_{\mathbb{R}^3} [uf(u) - \gamma F(u)] dx + \frac{2_s^* - \gamma}{2_s^*[(s+t)-3]} |u|_{2_s^*}^{2_s^*} \\ &\geq \frac{(s+t)\gamma - (4s+2t)}{2[(s+t)\gamma-3]} \|u\|_{E_\lambda}^2. \end{aligned}$$

It follows from  $(f_1) - (f_2)$  and (2.5) that

$$\int_{\mathbb{R}^3} [(s+t)f(u)u - 3F(u)] dx \leq \frac{2s+2t-3}{4} \|u\|_{E_\lambda}^2 + C_1 \|u\|_{E_\lambda}^q.$$

From which, combining  $(2s+2t-3)V(x) \geq 2(\nabla V, x) \geq 0$  for all  $x \in \mathbb{R}^3$  in  $(V_5)$  and (2.3), we see that

$$\frac{2s+2t-3}{4} \|u\|_{E_\lambda}^2 \leq C_1 \|u\|_{E_\lambda}^q + S_s^{-\frac{2}{2_s^*}} \|u\|_{E_\lambda}^{2_s^*}, \quad \forall u \in \mathcal{M}_\lambda,$$

yielding that  $\|u\|_{E_\lambda} \geq C_2$  for some  $C_2 > 0$  independent of  $\lambda$ . So, we arrive at (3.8).

On the other hand, we begin verifying (3.9). Without loss of generality, we are assuming that  $0 \in \Omega$ . Because  $\Omega$  is an open subset of  $\mathbb{R}^3$ , it holds that  $B_{r_0}(0) \subset \Omega$  for some  $r_0 > 0$ . Given a constant  $\hat{r}_0 > 0$

which will be determined later, we choose a cutoff function  $\psi \in C_0^\infty(\mathbb{R}^3)$  in such a way that  $\psi(x) \equiv 1$  if  $|x| \leq \hat{r}_0$  and  $\psi(x) \equiv 0$  if  $|x| \geq 2\hat{r}_0$ . For all  $\varepsilon > 0$ , we define

$$(3.10) \quad u_\varepsilon(x) = \psi(x)U_\varepsilon(x), \quad \forall x \in \mathbb{R}^3,$$

where  $U_\varepsilon(x) = \varepsilon^{-\frac{3-2s}{2}}u^*\left(\frac{x}{\varepsilon}\right)$ ,  $u^*(x) = \frac{U\left(x/S_s^{\frac{1}{2s}}\right)}{\|U\|_{2s}^*}$  and  $U(x) = \frac{\kappa}{(\tau^2+|x|^2)^{\frac{3-2s}{2}}}$  with  $\kappa \neq 0$  and  $\tau > 0$ . Due to Lemma 3.1 and (3.8), there exists a  $\theta_\varepsilon > 0$  such that

$$0 < m_\lambda \leq \max_{\theta > 0} I_\lambda(u_\theta) = I_\lambda((u_\varepsilon)_{\theta_\varepsilon}).$$

Next, we shall prove that there exist two constants  $\theta_*$ ,  $\theta^* > 0$  such that  $\theta_* \leq \theta_\varepsilon \leq \theta^*$ . First, we claim that  $\theta_\varepsilon$  is bounded from below by a positive constant. Otherwise, there is a sequence  $\varepsilon_n \rightarrow 0$  such that  $\theta_{\varepsilon_n} \rightarrow 0$ . Then, we conclude that  $(u_{\varepsilon_n})_{\varepsilon_n} \rightarrow 0$  in  $E_\lambda$ . So, we have

$$0 < m_\lambda \leq I_\lambda((u_\varepsilon)_{\theta_\varepsilon}) \rightarrow I_\lambda(0) = 0,$$

a contradiction. Taking some similar calculations in the proof of Lemma 3.1, one has  $\lim_{\theta_\varepsilon \rightarrow +\infty} I_\lambda((u_\varepsilon)_{\theta_\varepsilon}) = -\infty$  which is absurd, too. Thus, we conclude the claim. Letting  $\hat{r}_0 = \frac{1}{2}\theta_*r_0$ , then

$$\int_{\mathbb{R}^3} V(\theta_\varepsilon^{-1}x)u_\varepsilon^2 dx = \int_{B_{\theta_\varepsilon r_0}(0)} V(\theta_\varepsilon^{-1}x)u_\varepsilon^2 dx + \int_{\mathbb{R}^3 \setminus B_{\theta_\varepsilon r_0}(0)} V(\theta_\varepsilon^{-1}x)u_\varepsilon^2 dx = 0$$

from where it follows that

$$\begin{aligned} I_\lambda((u_\varepsilon)_{\theta_\varepsilon}) &= \frac{\theta_\varepsilon^{4s+2t-3}}{2} \int_{\mathbb{R}^3} |(-\Delta)^{\frac{s}{2}}u_\varepsilon|^2 dx + \frac{\theta_\varepsilon^{4s+2t-3}}{4} \int_{\mathbb{R}^3} \phi_{u_\varepsilon}^t u_\varepsilon^2 dx \\ &\quad - \theta^{-3} \int_{\mathbb{R}^3} F(\theta^{s+t}u_\varepsilon) dx - \frac{\theta^{2s^*(s+t)-3}}{2_s^*} \int_{\mathbb{R}^3} |u_\varepsilon|^{2_s^*} dx. \end{aligned}$$

Clearly, the proof of (3.9) would be done if  $I_\lambda((u_\varepsilon)_{\theta_\varepsilon}) < \frac{s}{3}S_s^{\frac{3}{2s}}$  for some suitably small  $\varepsilon > 0$ . Let us adopt the useful estimates in Lemma 3.2 and apply  $(f_3)$  to reach

$$\begin{aligned} I_\lambda((u_\varepsilon)_{\theta_\varepsilon}) &\leq \left( \frac{\theta_\varepsilon^{4s+2t-3}}{2} - \frac{\theta_\varepsilon^{(s+t)2_s^*-3}}{2_s^*} \right) S_s^{\frac{3}{2s}} + O(\varepsilon^{3-2s}) + C|u_\varepsilon|_2^2 + C|u_\varepsilon|_{\frac{12}{3+2t}}^4 - C\hat{\mu}|u_\varepsilon|_p^p \\ &\leq \frac{s}{3}S_s^{\frac{3}{2s}} + O(\varepsilon^{3-2s}) + C|u_\varepsilon|_2^2 + C|u_\varepsilon|_{\frac{12}{3+2t}}^4 - C\hat{\mu}|u_\varepsilon|_p^p, \end{aligned}$$

where we have used the following inequality

$$\int_{\mathbb{R}^3} \phi_{u_\varepsilon}^t u_\varepsilon^2 dx \leq C \left( \int_{\mathbb{R}^3} |u_\varepsilon|^{\frac{12}{3+2t}} dx \right)^{\frac{3+2t}{3}}.$$

To continue the proof, we divide the following three different cases.

**Case 1.**  $2 < \frac{3}{3-2s}$  which is equivalent to  $s > \frac{3}{4}$ . Then,

$$I_\lambda((u_\varepsilon)_{\theta_\varepsilon}) \leq \frac{s}{3}S_s^{\frac{3}{2s}} + O(\varepsilon^{3-2s}) + C|u_\varepsilon|_{\frac{12}{3+2t}}^4 - C\hat{\mu}|u_\varepsilon|_p^p.$$

**Case 2.**  $2 = \frac{3}{3-2s}$  which is equivalent to  $s = \frac{3}{4}$ . Then,

$$I_\lambda((u_\varepsilon)_{\theta_\varepsilon}) \leq \frac{s}{3}S_s^{\frac{3}{2s}} + O(\varepsilon^{2s}|\log \varepsilon|) + C|u_\varepsilon|_{\frac{12}{3+2t}}^4 - C\hat{\mu}|u_\varepsilon|_p^p.$$

**Case 3.**  $2 > \frac{3}{3-2s}$  which is equivalent to  $s < \frac{3}{4}$ . Then,

$$I_\lambda((u_\varepsilon)_{\theta_\varepsilon}) \leq \frac{s}{3}S_s^{\frac{3}{2s}} + O(\varepsilon^{2s}) + C|u_\varepsilon|_{\frac{12}{3+2t}}^4 - C\hat{\mu}|u_\varepsilon|_p^p.$$

We note that  $\frac{3s+t}{s+t} < \frac{2s}{3-2s} = \frac{3}{3-2s} - 1$  for any  $s > \frac{3}{4}$  and  $\frac{3s+t}{s+t} \geq \frac{2s}{3-2s} = \frac{3}{3-2s} - 1$  for any  $s \leq \frac{3}{4}$ . Thereby,

(a) If  $s > \frac{3}{4}$  in Case 1. It follows from (3.7) that

$$\lim_{\varepsilon \rightarrow 0^+} \frac{|u_\varepsilon|^4 \frac{12}{3+2t}}{\varepsilon^{3-2s}} \leq \begin{cases} \lim_{\varepsilon \rightarrow 0^+} \frac{O(\varepsilon^{4s+2t-3})}{\varepsilon^{3-2s}} = 0, & \frac{12}{3+2t} > \frac{3}{3-2s}, \\ \lim_{\varepsilon \rightarrow 0^+} \frac{O(\varepsilon^{4s+2t-3}) |\log \varepsilon|^{\frac{3+2t}{3}}}{\varepsilon^{3-2s}} = 0, & \frac{12}{3+2t} = \frac{3}{3-2s}, \\ \lim_{\varepsilon \rightarrow 0^+} \frac{O(\varepsilon^{2(3-2s)})}{\varepsilon^{3-2s}} = 0, & \frac{12}{3+2t} < \frac{3}{3-2s}. \end{cases}$$

Moreover, since  $\frac{4s}{3-2s} < p < \frac{6}{3-2s}$  gives that  $2s - \frac{3-2s}{2}p < 0$ , one infers from (3.7) again that

$$\lim_{\varepsilon \rightarrow 0^+} \frac{\hat{\mu} |u_\varepsilon|^p}{\varepsilon^{3-2s}} = \begin{cases} \lim_{\varepsilon \rightarrow 0^+} \hat{\mu} \frac{O(\varepsilon^{3-\frac{3-2s}{2}p})}{\varepsilon^{3-2s}} = +\infty, & \frac{4s}{3-2s} < p < \frac{6}{3-2s}, \\ \lim_{\varepsilon \rightarrow 0^+} \hat{\mu} \frac{O(\varepsilon^{3-\frac{3-2s}{2}p})}{\varepsilon^{3-2s}}, & \frac{3}{3-2s} < p \leq \frac{4s}{3-2s}, \\ \lim_{\varepsilon \rightarrow 0^+} \hat{\mu} \frac{O(\varepsilon^{3-\frac{3-2s}{2}p}) |\log \varepsilon|}{\varepsilon^{3-2s}}, & p = \frac{3}{3-2s}, \\ \lim_{\varepsilon \rightarrow 0^+} \hat{\mu} \frac{O(\varepsilon^{\frac{3-2s}{2}p})}{\varepsilon^{3-2s}}, & \frac{4s+2t}{s+t} < p < \frac{3}{3-2s}. \end{cases}$$

Choosing  $\hat{\mu} = \varepsilon^{-2s}$ , then the above three unknown limits would also be  $+\infty$ .

(b) If  $s = \frac{3}{4}$  in Case 2. Since  $\frac{12}{3+2t} > 2 - \frac{3}{3-2s}$ , there holds

$$\lim_{\varepsilon \rightarrow 0^+} \frac{|u_\varepsilon|^4 \frac{12}{3+2t}}{\varepsilon^{2s} |\log \varepsilon|} \leq \lim_{\varepsilon \rightarrow 0^+} \frac{O(\varepsilon^{4s+2t-3})}{\varepsilon^{2s} |\log \varepsilon|} = 0.$$

By  $\frac{3}{3-2s} = 2 < \frac{4s+2t}{s+t} < p$ , for any  $\hat{\mu} > 0$ , we have that

$$\lim_{\varepsilon \rightarrow 0^+} \frac{\hat{\mu} |u_\varepsilon|^p}{\varepsilon^{2s} |\log \varepsilon|} = \lim_{\varepsilon \rightarrow 0^+} \hat{\mu} \frac{O(\varepsilon^{3-\frac{3-2s}{2}p})}{\varepsilon^{2s} |\log \varepsilon|} = +\infty, \quad \frac{4s+2t}{s+t} < p < \frac{6}{3-2s}.$$

(c) If  $s < \frac{3}{4}$  in Case 3. Since  $\frac{3}{3-2s} \in (\frac{3}{2}, 2)$ , then  $\frac{12}{3+2t} > \frac{3}{3-2s}$  and  $\frac{3}{3-2s} < \frac{4s+2t}{s+t} < p < \frac{6}{3-2s}$ . Hence,

$$\lim_{\varepsilon \rightarrow 0^+} \frac{|u_\varepsilon|^4 \frac{12}{3+2t}}{\varepsilon^{2s}} \leq \lim_{\varepsilon \rightarrow 0^+} \frac{O(\varepsilon^{4s+2t-3})}{\varepsilon^{2s}} = 0$$

and for any  $\hat{\mu} > 0$ , there holds

$$\lim_{\varepsilon \rightarrow 0^+} \frac{\hat{\mu} |u_\varepsilon|^p}{\varepsilon^{2s}} = \lim_{\varepsilon \rightarrow 0^+} \hat{\mu} \frac{O(\varepsilon^{3-\frac{3-2s}{2}p})}{\varepsilon^{2s}} = +\infty, \quad \frac{4s+2t}{s+t} < p < \frac{6}{3-2s}.$$

At this stage, we will apply Point-(a) to Case 1, from (I) and (II) in (1.6); Point-(b) to Case 2 and Point-(c) to Case 3, from (III) in (1.6); there exists a sufficiently small  $\varepsilon > 0$  to arrive at the desired result. The proof is completed.  $\square$

As a byproduct of Lemma 3.3, we conclude that  $m_\lambda$  is well-defined. Before looking for a minimizer for it, we shall derive the following result which permits us to show that the weak limit of the minimizing sequence of  $m_\lambda$  is nontrivial.

**Lemma 3.4.** *Let  $s, t \in (0, 1)$  satisfy  $2s + 2t > 3$ . Assume that  $(V_1) - (V_5)$  and  $(f_1) - (f_5)$ . Let  $\lambda > \Lambda_0$  and  $(u_n) \subset E_\lambda$  be a minimizing sequence sequence of  $m_\lambda$ , then there exist  $r \in \left(2, \frac{3(3-s)}{3-2s}\right)$  and  $\sigma_0 > 0$ , independent of  $\lambda$ , such that  $|u_n|_r \geq \sigma_0$ , for all  $n \geq 1$ .*

*Proof.* First of all, we can show that  $(u_n)$  is uniformly bounded in  $n \in \mathbb{N}$  for all  $\lambda > \Lambda$ , see e.g. Lemma 3.5 below in detail. Let us divide the proof into intermediate steps.

STEP I: Let  $\lambda > \Lambda_0$  and  $(u_n) \subset E_\lambda$  be a minimizing sequence of  $m_\lambda$ , then there exist  $r \in \left(2, \frac{3(3-s)}{3-2s}\right)$  and  $\sigma = \sigma(\lambda) > 0$  such that  $|u_n|_r \geq \sigma$ , for all  $n \geq 1$ .

Suppose, by contradiction, that  $u_n \rightarrow 0$  in  $L^r(\mathbb{R}^3)$  for each  $r \in \left(2, \frac{3(3-s)}{3-2s}\right)$ . Due to the boundedness of  $(u_n)$  in  $E_\lambda$ , we see that  $(u_n)$  is uniformly bounded in  $L^q(\mathbb{R}^2)$  for all  $q \in (2, 2_s^*)$ , too. As a consequence, one simply arrives at

$$(3.11) \quad \lim_{n \rightarrow \infty} \int_{\mathbb{R}^3} \phi_{u_n}^t u_n^2 dx = 0, \quad \lim_{n \rightarrow \infty} \int_{\mathbb{R}^3} f(u_n) u_n dx = 0 \quad \text{and} \quad \lim_{n \rightarrow \infty} \int_{\mathbb{R}^3} F(u_n) dx = 0.$$

Without loss of generality, we could assume that  $\|u_n\|_{E_\lambda}^2 \rightarrow l$  as  $n \rightarrow \infty$ . Obviously, we derive  $l > 0$ . Otherwise,  $\|u_n\|_{E_\lambda}^2 \rightarrow 0$  and hence  $|u_n|_{2_s^*} \rightarrow 0$  as  $n \rightarrow \infty$  by (2.3). Combining these facts and (3.11), it holds that  $m_\lambda = \lim_{n \rightarrow \infty} I_\lambda(u_n) = 0$ , which is absurd because of (3.8). Now, we claim that  $\lim_{n \rightarrow \infty} |u_n|_{2_s^*}^2 = l$ . Indeed, according to  $G_\lambda(u_n) = 0$ , (3.11) and  $\frac{4s+2t-3}{2} = \frac{2_s^*(s+t)-3}{2_s^*}$  with  $(V_5)$ , we obtain the desired result. Using (2.3) again, then  $l \leq S_s^{-\frac{2_s^*}{2}} l^{\frac{2_s^*}{2}}$  which gives that  $l \geq S_s^{\frac{3}{2_s^*}}$ . So, it follows from (3.11) that

$$m_\lambda = \lim_{n \rightarrow \infty} I_\lambda(u_n) = \left(\frac{1}{2} - \frac{1}{2_s^*}\right) l \geq \frac{s}{3} S_s^{\frac{3}{2_s^*}}$$

reaching a contradiction with (3.9).

STEP II: Conclusion.

Let  $r \in \left(2, \frac{3(3-s)}{3-2s}\right)$  be as in Step I. Suppose by contradiction that the uniform control from below of  $L^r(\mathbb{R}^2)$ -norm is false. Then, for every  $k \in \mathbb{N}$ ,  $k \neq 0$ , there exist  $\lambda_k > \Lambda_0$  and a minimizing sequence  $(u_{k,n})$  of  $m_{\lambda_k}$  such that

$$|u_{k,n}|_r < \frac{1}{k}, \quad \text{definitely.}$$

Then, by a diagonalization argument, for any  $k \geq 1$ , we can find an increasing sequence  $(n_k)$  in  $\mathbb{N}$  and  $u_{n_k} \in E_{\lambda_{n_k}}$  such that

$$u_{n_k} \in \mathcal{M}_{\lambda_k}, \quad J_{n_k}(u_{n_k}) = c_{\lambda_{n_k}} + o_k(1), \quad |u_{n_k}|_r = o_k(1),$$

where  $o_k(1)$  is a positive quantity which goes to zero as  $k \rightarrow +\infty$ . Then, we are able to arrive at a same contradiction in the Step I with (3.9), again. The proof is completed  $\square$

**Lemma 3.5.** *Let  $s, t \in (0, 1)$  satisfy  $2s + 2t > 3$ . Assume that  $(V_1) - (V_5)$  and  $(f_1) - (f_5)$  with one of the assumptions in (1.6), then there is a  $\Lambda > 0$  such that  $m_\lambda$  can be attained for all  $\lambda > \Lambda$ .*

*Proof.* Let  $(u_n) \subset \mathcal{M}_\lambda$  be a sequence satisfying  $I_\lambda(u_n) \rightarrow m_\lambda$  as  $n \rightarrow \infty$ . First of all, we claim that  $(u_n)$  is uniformly bounded in  $E_\lambda$  with respect to  $n \in \mathbb{N}$  for all  $\lambda > \Lambda_0$ . Indeed, since  $(u_n) \subset \mathcal{M}_\lambda$  gives that  $G_\lambda(u_n) = 0$  and so

$$\begin{aligned} m_\lambda &= I_\lambda(u_n) + o_n(1) = I_\lambda(u_n) - \frac{1}{(s+t)\gamma - 3} G_\lambda(u_n) + o_n(1) \\ &= \frac{(s+t)\gamma - (4s+2t)}{2[(s+t)\gamma - 3]} |(-\Delta)^{\frac{s}{2}} u_n|_2^2 + \frac{1}{2[(s+t)\gamma - 3]} \int_{\mathbb{R}^3} \lambda [(s+t)(\gamma - 2)V(x) + (\nabla V, x)] u_n^2 dx \end{aligned}$$

$$\begin{aligned}
 & + \frac{(s+t)\gamma - (4s+2t)}{4[(s+t)\gamma - 3]} \int_{\mathbb{R}^3} \phi_{u_n}^t u_n^2 dx + \frac{s+t}{(s+t)\gamma - 3} \int_{\mathbb{R}^3} [u_n f(u_n) - \gamma F(u_n)] dx + o_n(1) \\
 & + \frac{2_s^* - \gamma}{2_s^*[(s+t) - 3]} \int_{\mathbb{R}^3} |u_n|^{2_s^*} dx + o_n(1) \\
 (3.12) \quad & \geq \frac{(s+t)\gamma - (4s+2t)}{2[(s+t)\gamma - 3]} |(-\Delta)^{\frac{s}{2}} u_n|_2^2 + \frac{(s+t)\gamma - (4s+2t)}{4[(s+t)\gamma - 3]} \int_{\mathbb{R}^3} \phi_{u_n}^t u_n^2 dx + o_n(1)
 \end{aligned}$$

which together with (3.9) implies that  $|(-\Delta)^{\frac{s}{2}} u_n|_2$  is uniformly bounded in  $n \in \mathbb{N}$  for all  $\lambda > \Lambda_0$ . By means of the interpolation inequality, for  $q \in (2, 2_s^*)$ , we combine (2.3) and (2.5) to derive

$$\begin{aligned}
 |u_n|_q^q & \leq |u_n|_2^{2\nu} |u_n|_{2_s^*}^{2_s^*(1-\nu)} \leq C \|u\|_{E_\lambda}^{2\nu} |u_n|_{2_s^*}^{2(1-\nu)} \\
 (3.13) \quad & \leq C \|u_n\|_{E_\lambda}^{2\nu} |(-\Delta)^{\frac{s}{2}} u_n|_2^{1-\nu} \leq C \|u_n\|_{E_\lambda}^{2\nu},
 \end{aligned}$$

where  $\nu = \frac{2_s^* - q}{2_s^* - 2} \in (0, 1)$ . Therefore, using  $(f_1) - (f_2)$ , it follows from (3.13), (2.3) and (3.9) that

$$\begin{aligned}
 m_\lambda & = I_\lambda(u_n) + o_n(1) \geq \frac{1}{4} \|u_n\|_{E_\lambda}^2 - C |u_n|_q^q - C |(-\Delta)^{\frac{s}{2}} u_n|_2^{2_s^*} \\
 & \geq \frac{1}{4} \|u_n\|_{E_\lambda}^2 - C \|u_n\|_{E_\lambda}^{2\nu} - C,
 \end{aligned}$$

yielding that  $(u_n)$  is uniformly bounded in  $E_\lambda$  with respect to  $n \in \mathbb{N}$  for all  $\lambda > \Lambda_0$  since  $\xi \in (0, 1)$ . So, up to a subsequence if necessary, there is a  $u \in E_\lambda$  such that  $u_n \rightharpoonup u$  in  $E_\lambda$ ,  $u_n \rightarrow u$  in  $L_{\text{loc}}^p(\mathbb{R}^3)$  for all  $2 < p < 2_s^*$  and  $u_n \rightarrow u$  a.e. in  $\mathbb{R}^3$ .

Secondly, we shall find a suitable large  $\Lambda > 0$  such that  $u \neq 0$  for all  $\lambda > \Lambda$ . Owing to the above discussions, we know that  $\|u_n\|_{E_\lambda}^2 \leq C^*$  for a suitable  $C^* > 0$ , for any  $n \geq 1$  and  $\lambda > \Lambda_0$ . Let  $r > 2$  and  $\sigma_0 > 0$  be given as in Lemma 3.4, recalling  $(V_3)$ , there is a sufficiently large constant  $\bar{R} > 1$  such that,

$$(3.14) \quad \int_{B_{\bar{R}}^c(0) \cap \Sigma} |u_n|^r dx \leq \frac{\sigma_0}{4}, \quad \text{for all } \lambda > \Lambda_0 \text{ and for all } n \geq 1.$$

Since  $V(x) \geq c$  on  $\Sigma^c$  by  $(V_3)$ , we have

$$\int_{B_{\bar{R}}^c(0) \cap \Sigma^c} |u_n|^2 dx \leq \frac{1}{\lambda c} \int_{B_{\bar{R}}^c(0) \cap \Sigma^c} \lambda V(x) |u_n|^2 dx \leq \frac{C^*}{\lambda c}$$

It easily infers that

$$\int_{B_{\bar{R}}^c(0) \cap \Sigma^c} |u_n|^r dx \leq \left( \int_{B_{\bar{R}}^c(0) \cap \Sigma^c} |u_n|^2 dx \right)^{\frac{1}{2}} \left( \int_{B_{\bar{R}}^c(0) \cap \Sigma^c} |u_n|^{2(r-1)} dx \right)^{\frac{1}{2}},$$

and so one can find a  $\Lambda > \Lambda_0$  such that

$$(3.15) \quad \int_{B_{\bar{R}}^c(0) \cap \Sigma^c} |u_n|^r dx \leq \frac{\sigma_0}{4}, \quad \text{for all } \lambda > \Lambda \text{ and for all } n \geq 1.$$

Finally, we fix  $\lambda > \Lambda_0$ , if  $u_n \rightharpoonup u \equiv 0$ , we can deduce that

$$(3.16) \quad \int_{B_{\bar{R}}(0)} |u_n|^r dx \leq \frac{\sigma_0}{4}, \quad \text{for all } n \text{ sufficiently large.}$$

Clearly, (3.14), (3.15) and (3.16) are in contradictions with Lemma 3.4.

Finally, we conclude that  $u_n \rightarrow u$  along a subsequence as  $n \rightarrow \infty$  for all  $\lambda > \Lambda$ . Define  $w_n \triangleq u_n - u$ , then thanks to Lemma 2.2-(3) and the Brézis-Lieb lemma,

$$(3.17) \quad \lim_{n \rightarrow \infty} I_\lambda(w_n) = \lim_{n \rightarrow \infty} [I_\lambda(u_n) - I_\lambda(u)] = m_\lambda - I_\lambda(u)$$

and

$$(3.18) \quad \lim_{n \rightarrow \infty} G_\lambda(w_n) = \lim_{n \rightarrow \infty} [G_\lambda(u_n) - G_\lambda(u)] = -G_\lambda(u).$$

We claim that  $G_\lambda(u) \leq 0$ . Otherwise, it has that  $\lim_{n \rightarrow \infty} G_\lambda(w_n) < 0$  by (3.18). Without loss of generality, we are assuming that  $G_\lambda(w_n) < 0$  for all  $n \in \mathbb{N}$ . From which, one knows that  $w_n \neq 0$  and so Lemma 3.1 permits us to determine a  $\theta_n > 0$  such that  $G_\lambda((w_n)_{\theta_n}) = 0$ . Combining (3.4) and (3.17)-(3.18),

$$\begin{aligned} m_\lambda - I_\lambda(u) + \frac{1}{4s+2t-3}G_\lambda(u) &= \lim_{n \rightarrow \infty} \left[ I_\lambda(w_n) - \frac{1}{4s+2t-3}G_\lambda(w_n) \right] \\ &\geq \lim_{n \rightarrow \infty} \left[ I_\lambda((w_n)_{\theta_n}) - \frac{\theta_n^{4s+2t-3}}{4s+2t-3}G_\lambda(w_n) \right] > \lim_{n \rightarrow \infty} I_\lambda((w_n)_{\theta_n}) \geq m_\lambda, \end{aligned}$$

which gives that

$$I_\lambda(u) - \frac{1}{4s+2t-3}G_\lambda(u) < 0.$$

It is similar to (3.12) that we would get a contradiction. Hence, we have arrived at  $G_\lambda(u) \leq 0$ . Adopting Lemma 3.1 again, there exists a  $\theta > 0$  such that  $u_\theta \in \mathcal{M}_\lambda$ . Owing to (3.4) and Fatou's lemma,

$$\begin{aligned} m_\lambda &= \lim_{n \rightarrow \infty} I_\lambda(u_n) = \lim_{n \rightarrow \infty} \left[ I_\lambda(u_n) - \frac{1}{4s+2t-3}G_\lambda(u_n) \right] \geq I_\lambda(u) - \frac{1}{4s+2t-3}G_\lambda(u) \\ &\geq I_\lambda(u_\theta) - \frac{\theta^{4s+2t-3}}{4s+2t-3}G_\lambda(u) \geq I_\lambda(u_\theta) \geq m_\lambda, \end{aligned}$$

which yields that  $u_n \rightarrow u$  in  $E_\lambda$ . Consequently,  $I_\lambda(u) = m_\lambda$  and  $G_\lambda(u) = 0$ . The proof is completed.  $\square$

#### 4. PROOF OF MAIN THEOREMS

**4.1. Proof of Theorem 1.1.** Now, we are in position to show the proof of Theorem 1.1.

The proof would be done if  $u$  obtained in Lemma 3.5 satisfies  $I'_\lambda(u) = 0$  in  $E_\lambda^{-1}$ . Motivated by [30], we argue it indirectly. If  $I'_\lambda(u) \neq 0$ , there exists a  $\varphi \in C_0^\infty(\mathbb{R}^3)$  such that  $I'_\lambda(u)\varphi < -1$ . Let  $\varepsilon > 0$  be small enough and satisfy

$$(4.1) \quad I'_\lambda(u_\theta + \tau\varphi)\varphi \leq -\frac{1}{2}, \text{ for } |\theta - 1| + |\tau| \leq \varepsilon.$$

Let  $\chi \in C_0^\infty(\mathbb{R}, [0, 1])$  be a cut-off function satisfying  $\chi(\theta) \equiv 1$  for every  $|\theta - 1| \leq \frac{\varepsilon}{2}$  and  $\chi(\theta) \equiv 0$  for all  $|\theta - 1| \geq \varepsilon$ . For any  $\theta > 0$ , we define

$$\eta(\theta) \triangleq \begin{cases} u_\theta, & \text{if } |\theta - 1| \geq \varepsilon, \\ u_\theta + \varepsilon\chi(\theta)\varphi, & \text{if } |\theta - 1| < \varepsilon. \end{cases}$$

Obviously,  $\eta \in \mathcal{C}(E_\lambda)$  and one can fix  $\varepsilon > 0$  sufficiently small such that  $\|\eta(\theta)\|_{E_\lambda} > 0$  for  $|\theta - 1| < \varepsilon$ . By (4.1), it is easy to show that

$$\max_{\theta > 0} I_\lambda(\eta(\theta)) < m_\lambda.$$

Proceeding as the proof of Lemma 3.1, we have  $G_\lambda(\eta(1 - \varepsilon)) > 0$  and  $G_\lambda(\eta(1 + \varepsilon)) < 0$ . Since  $G_\lambda(\eta(\theta))$  is continuous, there exists  $\theta_0 \in (1 - \varepsilon, 1 + \varepsilon)$  such that  $G_\lambda(\eta(\theta_0)) = 0$  which is  $\eta(\theta_0) \in \mathcal{M}_\lambda$ . Therefore,  $m_\lambda \leq I_\lambda(\eta(\theta_0)) \leq \max_{\theta > 0} I_\lambda(\eta(\theta)) < m_\lambda$ , a contradiction. As to the positivity of  $u$ , it is standard and we omit it here. The proof is completed.



Next, we will deal with the concentrating behaviour of ground state solutions obtained in Theorem 1.1. For any  $u \in H_0^s(\Omega)$ , we denote by  $\tilde{u} \in H^s(\mathbb{R}^3)$  its trivial extension, namely

$$\tilde{u} \triangleq \begin{cases} u & \text{in } \Omega, \\ 0 & \text{in } \Omega^c = \{x : x \in \mathbb{R}^3 \setminus \Omega\}. \end{cases}$$

We now define  $I_0|_\Omega : H_0^s(\Omega) \rightarrow \mathbb{R}$  as

$$I_0|_\Omega(u) = \frac{1}{2} \int_\Omega |(-\Delta)^{\frac{s}{2}} u|^2 dx + \frac{c_t}{4} \int_\Omega \int_\Omega \frac{u^2(x)u^2(y)}{|x-y|^{3-2t}} dx dy - \int_\Omega f(u)u dx - \frac{1}{2_s^*} \int_\Omega |u|^{2_s^*} dx$$

and consider the minimization problem

$$m_0|_\Omega \triangleq \inf_{u \in \mathcal{M}_0|_\Omega} I_0|_\Omega(u)$$

where

$$\mathcal{M}_0|_\Omega = \{u \in H_0^s(\Omega) \setminus \{0\} : G_0|_\Omega(u) = 0\}$$

denotes the corresponding manifold and  $G_0|_\Omega : H_0^s(\Omega) \rightarrow \mathbb{R}$  is given by

$$\begin{aligned} G_0|_\Omega(u) &= \frac{4s+2t-3}{2} \int_\Omega |(-\Delta)^{\frac{s}{2}} u|^2 dx + \frac{4s+2t-3}{4} c_t \int_\Omega \int_\Omega \frac{u^2(x)u^2(y)}{|x-y|^{3-2t}} dx dy \\ &\quad - \int_\Omega [(s+t)f(u)u - 3F(u)] dx - \frac{2_s^*(s+t)-3}{2_s^*} \int_\Omega |u|^{2_s^*} dx. \end{aligned}$$

We note that, up to the above trivial extension, there holds that  $\mathcal{M}_0|_\Omega \subset \mathcal{M}_\lambda$  for all  $\lambda > 0$ .

For each  $\lambda > \Lambda_0$ , we denote by  $u_\lambda \in E_\lambda$  a ground state solution of system (1.1), that is,  $I'_\lambda(u_\lambda) = 0$  and  $I_\lambda(u_\lambda) = m_\lambda$ . Then, we prove Theorem 1.3 as follows.

**4.2. Proof of Theorem 1.3.** Let  $\lambda_n \rightarrow +\infty$  as  $n \rightarrow +\infty$  and  $(u_{\lambda_n}) \subset E_{\lambda_n}$  be a sequence of ground state solutions of system (1.1), that is,  $I'_{\lambda_n}(u_{\lambda_n}) = 0$  and  $I_{\lambda_n}(u_{\lambda_n}) = m_{\lambda_n}$ . Up to a subsequence if necessary, by (3.8) and  $\mathcal{M}_0|_\Omega \subset \mathcal{M}_\lambda$ , for all  $\lambda > 0$ ,

$$(4.2) \quad 0 < \rho \leq \lim_{n \rightarrow \infty} I_{\lambda_n}(u_{\lambda_n}) \triangleq \tilde{m}_\Omega \leq m_0|_\Omega < +\infty.$$

Clearly,  $(u_{\lambda_n})$  is bounded in  $H^s(\mathbb{R}^3)$ . Thereby, up to a subsequence if necessary, there is a  $u_0 \in H^s(\mathbb{R}^3)$  such that  $u_{\lambda_n} \rightharpoonup u_0$  in  $H^s(\mathbb{R}^3)$  and  $u_{\lambda_n} \rightarrow u_0$  a.e. in  $\mathbb{R}^3$ . By means of Lemma 2.2-(3), we conclude that  $I_0|'_\Omega(u_0) = 0$ . We claim that  $u \equiv 0$  in  $\Omega^c$ . Otherwise, there is a compact subset  $\Theta_{u_0} \subset \Omega^c$  with  $\text{dist}(\Theta_{u_0}, \partial\Omega^c) > 0$  such that  $u_0 \neq 0$  on  $\Theta_{u_0}$  and by Fatou's lemma

$$(4.3) \quad \liminf_{n \rightarrow \infty} \int_{\mathbb{R}^3} u_n^2 dx \geq \int_{\Theta_{u_0}} u_0^2 dx > 0.$$

Moreover, there exists  $\varepsilon_0 > 0$  such that  $V(x) \geq \varepsilon_0$  for any  $x \in \Theta_{u_0}$  by the assumptions  $(V_1)$  and  $(V_2)$ . Combining  $(f_4)$  with  $\gamma > 2$  and (4.2)-(4.3), we reach

$$\begin{aligned} c_\Omega &\geq \liminf_{n \rightarrow \infty} \left\{ \frac{\gamma-2}{2\gamma} \int_{\mathbb{R}^2} \lambda_n V(x) u_n^2 dx - \frac{|\gamma-4|}{4\gamma} \int_{\mathbb{R}^3} \phi_{u_{\lambda_n}}^t u_{\lambda_n}^2 - \frac{2_s^* - \gamma}{2_s^* \gamma} \int_{\mathbb{R}^3} |u_{\lambda_n}|^{2_s^*} dx \right\} \\ &\geq \frac{(q-2)\varepsilon_0}{2q} \int_{\Theta_u} u_0^2 dx \liminf_{n \rightarrow \infty} \lambda_n - \hat{C} = +\infty, \end{aligned}$$

a contradiction, where  $\hat{C} > 0$  is independent of  $n \in \mathbb{N}$ . Therefore,  $u_0 \in H_0^s(\Omega)$  by the fact that  $\partial\Omega$  is smooth and  $I_0|'_\Omega(u_0) = 0$ . Similar to the proof of Lemma 3.5, one knows  $u_0 \neq 0$ . Proceeding as the proof of Lemma 2.1, it holds that  $G_0|_\Omega(u_0) = 0$ . In view of (4.2), by  $u_0 \in H_0^s(\Omega)$ , we use the Fatou's lemma to obtain

$$m_0|_\Omega \geq \tilde{m}_\Omega = \liminf_{n \rightarrow \infty} \left[ I_{\lambda_n}(u_{\lambda_n}) - \frac{1}{4s+2t-3} G_{\lambda_n}(u_{\lambda_n}) \right]$$

$$\geq I_0|_\Omega(u_0) - \frac{1}{4s+2t-3}G_0|_\Omega(u_0) = I_0|_\Omega(u_0) \geq m_0|_\Omega$$

yielding that  $u_{\lambda_n} \rightarrow u_0$  in  $H^s(\mathbb{R}^3)$  and  $I_0|_\Omega(u_0) = m_0|_\Omega$ . The proof is finished.

**4.3. Proof of Theorem 1.4.** In this section, we are going to contemplate the existence of positive solutions for system 1.1 with a wider class of  $V$  and  $f$ . Without  $(V_5)$  and  $(f_5)$ , one could not take advantage of the minimization constraint manifold method explored in Section 3. Whereas, because of  $(f_4)$ , it seems impossible to prove that the  $(PS)$  sequence is uniformly bounded. As a consequence, we shall depend on an indirect approach developed by Jeanjean [17].

**Proposition 4.1.** (See [17, Theorem 1.1 and Lemma 2.3]) *Let  $(X, \|\cdot\|)$  be a Banach space and  $T \subset \mathbb{R}^+$  be an interval, consider a family of  $C^1$  functionals on  $X$  of the form*

$$\Phi_\mu(u) = A(u) - \mu B(u), \quad \forall \mu \in T,$$

with  $B(u) \geq 0$  and either  $A(u) \rightarrow +\infty$  or  $B(u) \rightarrow +\infty$  as  $\|u\| \rightarrow +\infty$ . Assume that there exists two points  $v_1, v_2 \in X$  such that

$$c_\mu = \inf_{\gamma \in \Gamma} \sup_{\theta \in [0,1]} \Phi_\mu(\gamma(\theta)) > \max\{\Phi_\mu(v_1), \Phi_\mu(v_2)\}, \quad \forall \mu \in T,$$

where

$$\Gamma = \{\gamma \in C([0,1], X) : \gamma(0) = v_1, \gamma(1) = v_2\}.$$

Then, for almost every  $\mu \in T$ , there is a sequence  $(u_n(\mu)) \subset X$  such that

- (a)  $(u_n(\mu))$  is bounded in  $X$ ;
- (b)  $\Phi_\mu(u_n(\mu)) \rightarrow c_\mu$  and  $\Phi'_\mu(u_n(\mu)) \rightarrow 0$ ;
- (c) the map  $\mu \rightarrow c_\mu$  is non-increasing and left continuous.

Letting  $T = [\delta, 1]$ , where  $\delta \in (0, 1)$  is a positive constant. To apply Proposition 4.1, we will introduce a family of  $C^1$ -functionals on  $X = E_\lambda$  with the form

$$(4.4) \quad I_{\lambda,\mu}(u) = \frac{1}{2} \int_{\mathbb{R}^3} [(-\Delta)^{\frac{s}{2}} u]^2 + \lambda V(x)|u|^2 dx + \frac{1}{4} \int_{\mathbb{R}^3} \phi_u^t u^2 dx - \mu \int_{\mathbb{R}^3} G(u) dx,$$

where and in the sequel  $G(z) = F(z) + \frac{1}{2_s^*}|z|^{2_s^*}$  for all  $z \in \mathbb{R}$ . Define  $I_{\lambda,\mu}(u) = A(u) - \mu B(u)$ , where

$$A(u) = \frac{1}{2} \int_{\mathbb{R}^3} [(-\Delta)^{\frac{s}{2}} u]^2 + \lambda V(x)|u|^2 dx + \frac{1}{4} \int_{\mathbb{R}^3} \phi_u^t u^2 dx \rightarrow +\infty \text{ as } \|u\|_{E_\lambda} \rightarrow +\infty,$$

and

$$B(u) = \int_{\mathbb{R}^3} G(u) dx \geq 0.$$

Clearly,  $I_{\lambda,\mu}$  is of class  $C^1$ -functionals with

$$I'_{\lambda,\mu}(u)v = \int_{\mathbb{R}^3} [(-\Delta)^{\frac{s}{2}} u (-\Delta)^{\frac{s}{2}} v + \lambda V(x)uv] dx + \int_{\mathbb{R}^3} \phi_u^t uv dx - \mu \int_{\mathbb{R}^3} g(u)v dx$$

for all  $u, v \in E_\lambda$ , where  $g(z) = f(z) + |z|^{2_s^*-2}z$  for all  $z \in \mathbb{R}$ .

For simplicity, from now on until the end of this section, we shall always suppose the assumptions in Theorem 1.4 when there is no misunderstanding.

**Lemma 4.2.** *The functional  $I_{\lambda,\mu}$  possesses a mountain-pass geometry, that is,*

- (a) *there exists  $v \in E_\lambda \setminus \{0\}$  independent of  $\mu$  such that  $I_{\lambda,\mu}(v) \leq 0$  for all  $\mu \in [\delta, 1]$ ;*
- (b)  *$c_{\lambda,\mu} \triangleq \inf_{\eta \in \Gamma} \sup_{\theta \in [0,1]} I_{\lambda,\mu}(\eta(\theta)) > \max\{I_{\lambda,\mu}(0), I_{\lambda,\mu}(v)\}$  for all  $\mu \in [\delta, 1]$ , where*

$$\Gamma = \{\eta \in C([0,1], E_\lambda) : \eta(0) = 0, \eta(1) = v\}.$$

*Proof.* The proof is very similar to the calculations on finding the existence of critical points in the proof of Lemma 3.1, so we omit the details.  $\square$

Repeating the arguments explored in Lemma 3.3, there is a constant  $\hat{\rho} > 0$  such that

$$(4.5) \quad \hat{\rho} \leq \inf_{\lambda > \Lambda_0} c_{\lambda, \mu} \leq \sup_{\lambda > \Lambda_0} c_{\lambda, \mu} < \frac{s}{3\mu^{\frac{3-2s}{2s}}} S_s^{\frac{3}{2s}}, \quad \forall \mu \in [\delta, 1].$$

**Lemma 4.3.** *Let  $(u_n)$  be a bounded (PS) sequence of the functional  $I_{\lambda, \mu}$  at the level  $c > 0$ , then for each  $\hat{M} \in \left( c, \frac{s}{3\mu^{\frac{3-2s}{2s}}} S_s^{\frac{3}{2s}} \right)$ , there exists a  $\hat{\Lambda} = \Lambda(\hat{M}) > 0$  such that  $(u_n)$  contains a strongly convergent subsequence in  $E_\lambda$  for all  $\lambda > \hat{\Lambda}$ .*

*Proof.* Since  $(u_n)$  is bounded in  $E_\lambda$ , then there exists a  $u \in E_\lambda$  such that  $u_n \rightharpoonup u$  in  $E_\lambda$ ,  $u_n \rightarrow u$  in  $L_{\text{loc}}^p(\mathbb{R}^3)$  with  $p \in [1, 2_s^*)$  and  $u_n \rightarrow u$  a.e. in  $\mathbb{R}^3$ . To show the proof clearly, we shall split it into several steps:

**Step 1:**  $I'_{\lambda, \mu}(u) = 0$  and  $I_{\lambda, \mu}(u) \geq 0$ .

To show  $I'_\lambda(u) = 0$ , since  $C_0^\infty(\mathbb{R}^3)$  is dense in  $E_\lambda$ , then it suffices to exhibit that  $I'_{\lambda, \mu}(u)\varphi = 0$  for every  $\varphi \in C_0^\infty(\mathbb{R}^3)$ . Thanks to Lemma 2.2-(3), it is a direct conclusion. Because  $u$  is a critical point of  $I_{\lambda, \mu}$ , according to Lemma 2.1, there holds  $P_{\lambda, \mu}(u) \equiv 0$ , where

$$\begin{aligned} P_{\lambda, \mu}(u) \triangleq & \frac{3-2s}{2} \int_{\mathbb{R}^3} |(-\Delta)^{\frac{s}{2}} u|^2 dx + \frac{1}{2} \int_{\mathbb{R}^3} [3V(x) + (\nabla V, x)] |u|^2 dx + \frac{2t+3}{4} \int_{\mathbb{R}^3} \phi_u^t u^2 dx \\ & - 3\mu \int_{\mathbb{R}^3} F(u) dx - \frac{\mu}{2_s^*} \int_{\mathbb{R}^3} |u|^{2_s^*} dx. \end{aligned}$$

Moreover, one easily sees that  $I'_{\lambda, \mu}(u)u = 0$  and so

$$I_{\lambda, \mu}(u) = I_{\lambda, \mu}(u) - \frac{1}{(s+t)\gamma-3} [(s+t)I'_{\lambda, \mu}(u)u - P_{\lambda, \mu}(u)] \geq 0$$

proving the Step 1.

**Step 2:** Define  $v_n \triangleq u_n - u$ , then there exists a  $\hat{\Lambda} = \Lambda(\hat{M}) > 0$  such that  $v_n \rightarrow 0$  in  $L^q(\mathbb{R}^3)$  for all  $q \in (2, 2_s^*)$  along a subsequence as  $n \rightarrow \infty$  when  $\lambda > \hat{\Lambda}$ .

Actually, since  $(v_n)$  is uniformly bounded in  $n \in \mathbb{N}$  for all  $\lambda > \Lambda_0$ , then we have one of the following two possibilities for some  $r > 0$ :

$$\begin{cases} \text{(i)} & \limsup_{n \rightarrow \infty} \sup_{y \in \mathbb{R}^3} \int_{B_r(y)} |v_n|^2 dx > 0, \\ \text{(ii)} & \limsup_{n \rightarrow \infty} \sup_{y \in \mathbb{R}^3} \int_{B_r(y)} |v_n|^2 dx = 0. \end{cases}$$

As a consequence, the conclusion would be clear if we could demonstrate that the case (i) cannot occur for sufficiently large  $\lambda > 0$ . Now, we suppose, by contradiction, that (i) was true. Proceeding as the very similar way in Lemma 3.5, there is a constant  $\hat{\delta} > 0$  independent of  $\lambda > \Lambda_0$  such that

$$\limsup_{n \rightarrow \infty} \sup_{y \in \mathbb{R}^3} \int_{B_r(y)} |v_n|^2 dx \geq \hat{\delta}$$

for some  $r > 0$ . Since  $(u_n)$  is uniformly bounded in  $E_\lambda$ , without loss of generality, we can assume that  $\lim_{n \rightarrow \infty} \|u_n\|_{E_\lambda}^2 \leq \Theta$  for some  $\Theta \in (0, +\infty)$ . Clearly, there holds  $\lim_{n \rightarrow \infty} \|v_n\|_{E_\lambda}^2 \leq 4\Theta$ . Recalling  $v_n \rightarrow 0$  in  $L_{\text{loc}}^q(\mathbb{R}^3)$  with  $q \in (2, 2_s^*)$  and  $|\mathcal{A}_R| \rightarrow 0$  as  $R \rightarrow +\infty$  by  $(V_2)$ , where  $\mathcal{A}_R \triangleq \{x \in \mathbb{R}^3 \setminus B_R(0) : V(x) < c\}$ ,

we can determine a sufficiently large but fixed  $R > 0$  to satisfy

$$(4.6) \quad \limsup_{n \rightarrow \infty} \int_{B_R(0)} |v_n|^2 dx < \frac{\hat{\delta}}{4}$$

and

$$(4.7) \quad |\mathcal{A}_R| < \left( \frac{\hat{\delta} S_s}{16\Theta} \right)^{\frac{q}{q-2}} |\Sigma|^{-\frac{2(2_s^*-q)}{2_s^*(q-2)}}.$$

Combining (2.5) and (4.7), one sees that

$$(4.8) \quad \limsup_{n \rightarrow \infty} \int_{\mathcal{A}_R} |v_n|^2 dx \leq \limsup_{n \rightarrow \infty} \left( \int_{\mathcal{A}_R} |v_n|^q dx \right)^{\frac{2}{q}} |\mathcal{A}_R|^{\frac{q-2}{q}} \leq 4\Theta |\Sigma|^{\frac{2(2_s^*-q)}{2_s^*q}} S_s^{-1} |\mathcal{A}_R|^{\frac{q-2}{q}} < \frac{\hat{\delta}}{4}.$$

Let us choose  $\hat{\Lambda} = \max \left\{ 1, \Lambda_0, \frac{16\Theta}{\hat{\delta}c} \right\}$ , then for all  $\lambda > \hat{\Lambda}$ , we reach

$$(4.9) \quad \limsup_{n \rightarrow \infty} \int_{\mathcal{B}_R} |v_n|^2 dx \leq \limsup_{n \rightarrow \infty} \frac{1}{\lambda c} \int_{\mathcal{B}_R} \lambda V(x) |v_n|^2 dx \leq \frac{4\Theta}{\lambda c} < \frac{\hat{\delta}}{4},$$

where  $\mathcal{B}_R \triangleq \{x \in \mathbb{R}^3 \setminus B_R(0) : V(x) \geq c\}$ . We gather (4.6), (4.7) and (4.9) to derive

$$\begin{aligned} \hat{\delta} &\leq \limsup_{n \rightarrow \infty} \sup_{y \in \mathbb{R}^3} \int_{B_r(y)} |v_n|^2 dx \leq \limsup_{n \rightarrow \infty} \int_{\mathbb{R}^3} |v_n|^2 dx \\ &= \limsup_{n \rightarrow \infty} \left( \int_{\mathbb{R}^3 \setminus B_R(0)} |v_n|^2 dx + \int_{B_R(0)} |v_n|^2 dx \right) \leq \frac{3\hat{\delta}}{4} \end{aligned}$$

which is impossible. The proof of this step is done.

**Step 3:** Passing to a subsequence if necessary,  $u_n \rightarrow u$  in  $E_\lambda$  as  $n \rightarrow \infty$ .

Since  $v_n \triangleq u_n - u$ , by Lemma 2.2-(3) and the Brézis-Lieb lemma, one has

$$(4.10) \quad I_{\lambda,\mu}(v_n) = I_{\lambda,\mu}(u_n) - I_{\lambda,\mu}(u) + o_n(1) \text{ and } I'_{\lambda,\mu}(v_n) = I'_\lambda(u_n) + o_n(1).$$

According to Step 2, we take advantage of (2.7) and  $(f_1) - (f_2)$  to deduce that

$$\lim_{n \rightarrow \infty} \int_{\mathbb{R}^3} \phi_{v_n}^t v_n^2 dx = 0 \text{ and } \lim_{n \rightarrow \infty} \int_{\mathbb{R}^3} f(v_n) v_n dx = 0$$

jointly with Lemma 2.2-(3) and the Brézis-Lieb lemma indicate that

$$o_n(1) = I'_\lambda(u_n)(u_n - u) - I'_\lambda(u)(u_n - u) = \|v_n\|_{E_\lambda}^2 - \mu |v_n|_{2_s^*}^{2_s^*}.$$

Let us suppose that  $\|v_n\|_{E_\lambda}^2 \rightarrow l$  and  $\mu |v_n|_{2_s^*}^{2_s^*} \rightarrow l$  along some subsequences and so

$$(4.11) \quad c \geq c - I_\lambda(u) = \lim_{n \rightarrow \infty} I_\lambda(v_n) = \left( \frac{1}{2} - \frac{1}{2_s^*} \right) l,$$

where we have used the Step 1 and (4.10). In view of (2.3), it holds that

$$(4.12) \quad (\mu^{-1}l)^{\frac{2}{2_s^*}} \leq S_s^{-1}l.$$

If  $l \neq 0$ , that is,  $l > 0$ , then  $l \geq \mu^{-\frac{3-2s}{2s}} S_s^{\frac{3}{2s}}$  by (4.12). As a consequence, with the help of (4.11), we arrive at  $c \geq \frac{s}{3\mu^{\frac{3-2s}{2s}}} S_s^{\frac{3}{2s}}$ , a contradiction. Therefore,  $l = 0$  which is the desired result. The proof is completed.  $\square$

Let us recall Proposition 4.1, Lemma 4.2 and Lemma 4.3, there are two sequences  $(\mu_n) \subset [\delta, 1]$  and  $(u_n) \subset E_\lambda \setminus \{0\}$  such that

$$(4.13) \quad I'_{\lambda, \mu_n}(u_n) = 0, \quad I_{\lambda, \mu_n}(u_n) = c_{\lambda, \mu_n} \quad \text{and} \quad \mu_n \rightarrow 1^-.$$

With (4.13) in hands, we are able to derive the proof of Theorem 1.4.

**Proof of Theorem 1.4.** First of all, since  $I'_{\lambda, \mu_n}(u_n) = 0$ , we are derived from a similar argument in Lemma 2.1 that  $P_{\lambda, \mu_n}(u_n) \equiv 0$ , where

$$\begin{aligned} P_{\lambda, \mu_n}(u) &\triangleq \frac{3-2s}{2} \int_{\mathbb{R}^3} |(-\Delta)^{\frac{s}{2}} u_n|^2 dx + \frac{1}{2} \int_{\mathbb{R}^3} [3V(x) + (\nabla V, x)] |u_n|^2 dx + \frac{2t+3}{4} \int_{\mathbb{R}^3} \phi_{u_n}^t u_n^2 dx \\ &\quad - 3\mu_n \int_{\mathbb{R}^3} F(u_n) dx - \frac{\mu_n}{2_s^*} \int_{\mathbb{R}^3} |u|^{2_s^*} dx. \end{aligned}$$

Proceeding as the proof of Lemma 3.5, one sees that  $(u_n)$  is uniformly bounded in  $E_\lambda$  for all  $\lambda > \Lambda_0$ .

Then, we claim that  $(u_n)$  is a  $(PS)_{c_{\lambda, 1}}$  sequence of the functional  $I_\lambda = I_{\lambda, 1}$ . Actually, taking into account  $\mu_n \rightarrow 1^-$  and Lemma 4.1-(c),

$$\lim_{n \rightarrow \infty} I_{\lambda, 1}(u_n) = \left( \lim_{n \rightarrow \infty} I_{\lambda, \mu_n}(u_n) + (\mu_n - 1) \int_{\mathbb{R}^3} G(u_n) dx \right) = \lim_{n \rightarrow \infty} c_{\lambda, \mu_n} = c_{\lambda, 1},$$

where we have used the fact that  $(G(u_n))$  is uniformly bounded in  $L^1(\mathbb{R}^3)$ . Similarly,

$$\begin{aligned} \lim_{n \rightarrow \infty} \frac{|I'_{\lambda, 1}(u_n)\psi|}{\|\psi\|_{E_\lambda}} &= \lim_{n \rightarrow \infty} \frac{|I'_{\lambda, \mu_n}(u_n)\psi + (\mu_n - 1) \int_{\mathbb{R}^3} g(u_n)\psi dx|}{\|\psi\|_{E_\lambda}} \\ &\leq \lim_{n \rightarrow \infty} \frac{|\mu_n - 1| \left| \int_{\mathbb{R}^3} g(u_n)\psi dx \right|}{\|\psi\|_{E_\lambda}} = 0, \quad \forall \psi \in E_\lambda. \end{aligned}$$

As a consequence, one has that  $(u_n)$  is a  $(PS)_{c_{\lambda, 1}}$  sequence of the functional  $I_\lambda = I_{\lambda, 1}$ .

Finally, combining the above two steps and (4.5), we can apply Lemma 4.3 to finish the proof.  $\square$

## REFERENCES

- [1] C.O. Alves, O.H. Miyagaki, Existence and concentration of solutions for a class of fractional elliptic equation in  $\mathbb{R}^N$  via penalization method, *Calc. Var. Partial Differential Equations*, **55** (2016), 1–19. [3](#)
- [2] A. Ambrosetti, D. Ruiz, Multiple bound states for the Schrödinger-Poisson equation, *Commun. Contemp. Math.*, **10** (2008), 1–14. [2](#)
- [3] V. Ambrosio, Ground states solutions for a non-linear equation involving a pseudo-relativistic Schrödinger operator, *J. Math. Phys.*, **57** (2016), 051502. [3](#)
- [4] A. Azzollini, P. d’Avenia, A. Pomponio, On the Schrödinger-Maxwell equations under the effect of a general nonlinear term, *Ann. Inst. H. Poincaré. Anal. Non-linéaire*, **27** (2010), 779–791. [2](#)
- [5] T. Bartsch, A. Pankov, Z. Wang, Nonlinear Schrödinger equations with steep potential well, *Commun. Contemp. math.*, **3** (2001), 549–569. [2, 4](#)
- [6] T. Bartsch, Z. Wang, Existence and multiplicity results for superlinear elliptic problems on  $\mathbb{R}^N$ , *Comm. Partial Differential Equations*, **20** (1995), 1725–1741. [2](#)
- [7] J. Bellazzini, L. Jeanjean, On dipolar quantum gases in the unstable regime, *SIAM J. Math. Anal.*, **48** (2017), 2028–2058. [2](#)
- [8] V. Benci, D. Fortunato, An eigenvalue problem for the Schrödinger-Maxwell equations, *Topol. Methods. Nonlinear Anal.* **11** (1998) 283-293. [2](#)
- [9] V. Benci, D. Fortunato, Solitary waves of the nonlinear Klein-Gordon coupled with Maxwell equations, *Rev. Math. Phys.* **14** (2002) 409-420. [2](#)
- [10] G. Bisci, V.D. Rădulescu, Ground state solutions of scalar field fractional Schrödinger equations, *Calc. Var. Partial Differential Equations*, **54** (2015), 2985–3008. [3](#)
- [11] L. Caffarelli, L. Silvestre, An extension problem related to the fractional Laplacian, *Comm. Partial Differential Equations*, **32** (2007), 1245–1260. [3](#)
- [12] R. Carles, P. Markowich, C. Sparber, On the Gross-Pitaevskii equation for trapped dipolar quantum gases, *Nonlinearity*, **21** (2008), 2569–2590. [2](#)

- [13] A. Cotsiolis, N. Tavoularis, Best constants for Sobolev inequalities for higher order fractional derivatives, *J. Math. Anal. Appl.*, **295** (2004), 225–236. [6](#)
- [14] X. He, Y. Meng, M. Squassina, Normalized solutions for a fractional Schrödinger-Poisson system with critical growth, *Calc. Var. Partial Differential Equations*, (2024), 63:142. [3](#)
- [15] X. He, W. Zou, Existence and concentration result for the fractional Schrödinger equations with critical nonlinearities, *Calc. Var. Partial Differential Equations*, **55** (2016), 1–39. [2](#), [3](#)
- [16] L. Huang, E. Rocha, J. Chen, Two positive solutions of a class of Schrödinger-Poisson system with indefinite nonlinearity, *J. Differential Equations*, **255** (2013), 2463–2483. [2](#)
- [17] L. Jeanjean, On the existence of bounded Palais-Smale sequences and application to a Landesman-Lazer type problem set on  $\mathbb{R}^N$ , *Proc. Roy. Soc. Edinburgh Sect. A*, **129** (1999) 787–809. [2](#), [18](#)
- [18] Y. Jiang, H. Zhou, Schrödinger-Poisson system with steep potential well, *J. Differential Equations*, **251** (2011), 582–608. [2](#)
- [19] J. Lan, X. He, On a fractional Schrödinger-Poisson system with doubly critical growth and a steep potential well, *J. Geom. Anal.*, **33** (2023), no. 6, Paper No. 187, 41 pp. [3](#)
- [20] N. Laskin, Fractional quantum mechanics and Lévy path integrals, *Phys. Lett. A*, **268** (2000) 298–305. [3](#)
- [21] N. Laskin, Fractional Schrödinger equation, *Phys. Rev. E*, **66** (2002) 56–108. [3](#)
- [22] H. Liu, L. Zhao, Ground-state solution of a nonlinear fractional Schrödinger-Poisson system, *Math. Methods Appl. Sci.*, **45** (2022), no. 4, 1934–1958. [3](#)
- [23] W. Long, J. Yang, W. Yu, Nodal solutions for fractional Schrödinger-Poisson problems, *Sci. China Math.*, **63** (2020), no. 11, 2267–2286. [3](#)
- [24] P. Lushnikov, Collapse and stable self-trapping for Bose-Einstein condensates with  $1/r^b$  type attractive interatomic interaction potential, *Phys. Rev. A*, **82** (2010), 023615. [2](#)
- [25] E. Di Nezza, G. Palatucci, E. Valdinoci, Hitchhiker’s guide to the fractional Sobolev spaces, *Bull. Sci. Math.* **136** (2012) 521–573. [5](#), [6](#), [8](#)
- [26] A. Pomponio, L. Shen, X. Zeng, Y. Zhang, Generalized Chern-Simons-Schrödinger system with sign-changing steep potential well: critical and subcritical exponential case, *J. Geom. Anal.*, **33** (2023), no. 6, Paper No. 185, 34 pp. [5](#)
- [27] P.H. Rabinowitz, *Minimax Methods in Critical Point Theory with Applications to Differential Equations*, CBMS Reg. Conf. Ser. Math. **65**, Amer. Math. Soc. Providence, 1986. [2](#)
- [28] D. Ruiz, Semiclassical states for coupled Schrödinger-Maxwell equations concentration around a sphere, *Math. Models Methods Appl. Sci.*, **15** (2005), 141–164. [2](#)
- [29] D. Ruiz, The Schrödinger-Poisson equation under the effect of a nonlinear local term, *J. Funct. Anal.*, **237** (2006), 655–674. [2](#)
- [30] D. Ruiz, G. Siciliano, Existence of ground states for a modified nonlinear Schrödinger equation, *Nonlinearity*, **23** (2010), 1221–1233. [16](#)
- [31] S. Secchi, Ground state solutions for nonlinear fractional Schrödinger equations in  $\mathbb{R}^N$ , *J. Math. Phys.* **54** (2013) 031501. [9](#)
- [32] L. Shen, Existence result for fractional Schrödinger-Poisson systems involving a Bessel operator without Ambrosetti-Rabinowitz condition, *Comput. Math. Appl.*, **75** (2018), no. 1, 296–306. [3](#)
- [33] L. Shen, Multiplicity and concentration results for fractional Schrödinger-Poisson systems involving a Bessel operator, *Math. Methods Appl. Sci.*, **41** (2018), no. 17, 7599–7611. [3](#)
- [34] L. Shen, Multiplicity and concentration results for fractional Schrödinger system with steep potential wells, *J. Math. Anal. Appl.*, **475** (2019), no. 2, 1385–1403. [3](#)
- [35] L. Shen, M. Squassina, Planar Schrödinger-Poisson system with steep potential well: supercritical exponential case, arXiv:2401.10663. [5](#)
- [36] L. Shen, M. Squassina, X. Zeng, Infinitely many solutions for a class of fractional Schrödinger equations coupled with neutral scalar field, *Discrete Contin. Dyn. Syst. S*, (2024), <https://doi.org/10.3934/dcdss.2024084>. [3](#)
- [37] L. Shen, X. Yao, Least energy solutions for a class of fractional Schrödinger-Poisson systems, *J. Math. Phys.*, **59** (2018), no. 8, 081501, 21 pp. [3](#), [8](#)
- [38] J. Sun, S. Ma, Ground state solutions for some Schrödinger-Poisson systems with periodic potentials, *J. Differential Equations*, **260** (2016), 2119–2149. [2](#)
- [39] K. Teng, Existence of ground state solutions for the nonlinear fractional Schrödinger-Poisson system with critical Sobolev exponent, *J. Differential Equations*, **261** (2016), 3061–3106. [3](#), [8](#)
- [40] K. Teng, Corrigendum to “Existence of ground state solutions for the nonlinear fractional Schrödinger-Poisson system with critical Sobolev exponent” [*J. Differential Equations* 261 (6) (2016) 3061–3106], *J. Differential Equations*, **262** (2017), 3132–3138. [10](#)
- [41] K. Teng, Y. Cheng, Multiplicity and concentration of nontrivial solutions for fractional Schrödinger-Poisson system involving critical growth, *Nonlinear Anal.*, **202** (2021), Paper No. 112144, 32 pp. [3](#)
- [42] M. Willem, *Minimax Theorems*, Birkhäuser, Boston, 1996. [2](#)

- [43] L. Yang, Z. Liu, Multiplicity and concentration of solutions for fractional Schrödinger equation with sublinear perturbation and steep potential well, *Comput. Math. Appl.*, **72** (2016), 1629–1640. [3](#)
- [44] J. Zhang, J.M. do Ó, M. Squassina, Fractional Schrödinger-Poisson systems with a general subcritical or critical nonlinearity, *Adv. Nonlinear Stud.*, **16** (2016), 15–30. [3](#)
- [45] L. Zhao, H. Liu, F. Zhao, Existence and concentration of solutions for Schrödinger-Poisson equations with steep well potential, *J. Differential Equations*, **255** (2013), 1–23. [2](#), [3](#), [4](#)
- [46] L. Zhao, F. Zhao, On the existence of solutions for the Schrödinger-Poisson equations, *J. Math. Anal. Appl.*, **346** (2008), 155–169. [2](#)

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