

- Lectures on DIFFERENTIAL GEOMETRY AND TOPOLOGY

Lecture XLIX

Prof. Mauro Spora - Dipartimento di
Matematica e Fisica "N. Tanzi"
UCSC - Brescia

MAYER-VIETORIS CALCULATIONS
(continued)

★ Further MV - Calculations

Σ_g : ^{connected} closed orientable surface (connected sum of g tori)
" ^{compact} + no boundary
 $g=0$: Sphere

Let us prove that $\chi(\Sigma_g) = 2 - 2g$

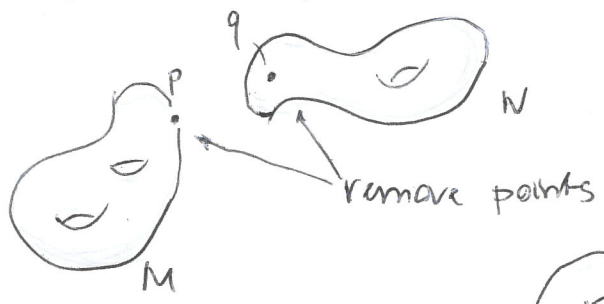
(we have already shown this for $g=0, 1$, and we hinted at the general case via "the canonical dissection").

Given that $H^0(\Sigma_g) \cong H^2(\Sigma_g) \cong \mathbb{R}$ (e.g. by Poincaré duality for compact oriented manifolds), it would easily follow that

$H^1(\Sigma_g) \cong \mathbb{R}^{2g}$

Given manifolds M, N , recall

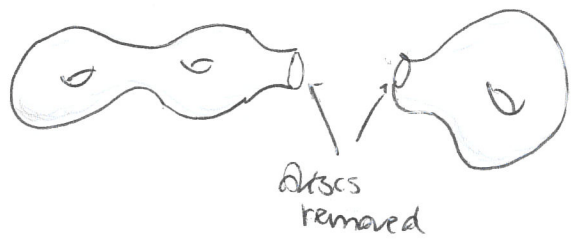
how to define their connected sum:



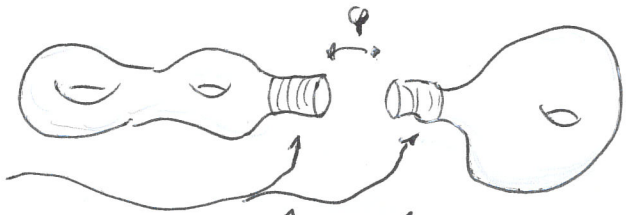
$U := M - \{p\} \cong M - B_n$

$V := N - \{q\} \cong N - B_n$

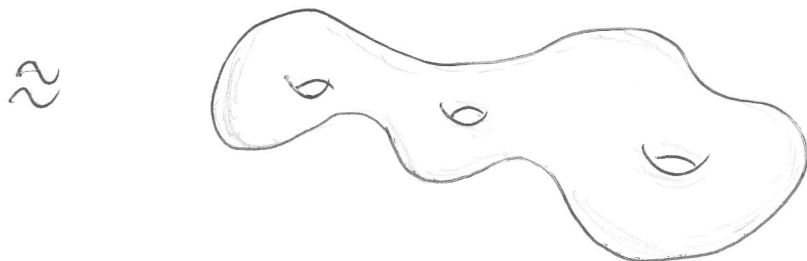
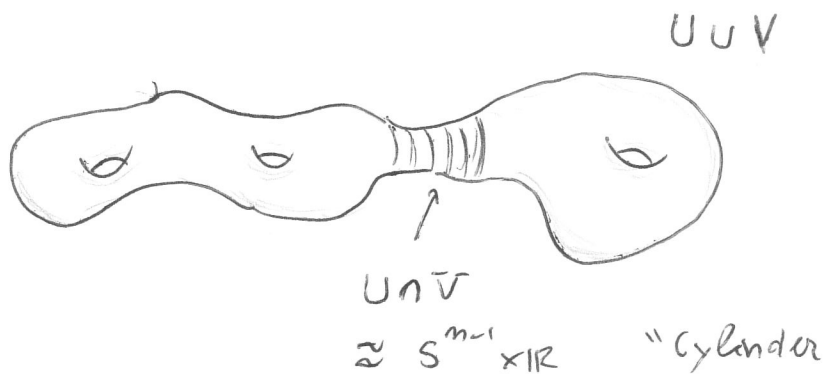
↑
open n -dim ball



glue along a collar $\cong S^{n-1} \times \mathbb{R}$



↑
"glue"
via a homeomorphism
(diffeomorphism)



(up to homeomorphisms (actually diffeomorphisms),
the procedure is independent of all choices..)

In our case ($\dim = 2$) we have

$$U = M - \{p\}, \quad \bar{V} = N - \{q\} \quad U \cup \bar{V} \approx M \# N$$

$$U \cap \bar{V} \approx S^1 \times \mathbb{R} \quad \text{cylinder} \quad \text{[cylinder icon]}$$

After these preparations, consider the MV - sequence
(long exact)

$$0 \rightarrow H^0(M \# N) \rightarrow H^0(U) \oplus H^0(\bar{V}) \rightarrow H^0(U \cap \bar{V}) \xrightarrow{d^*}$$

$$\hookrightarrow H^1(M \# N) \rightarrow H^1(U) \oplus H^1(\bar{V}) \rightarrow H^1(U \cap \bar{V}) \xrightarrow{d^*}$$

$$\hookrightarrow H^2(M \# N) \rightarrow H^2(U) \oplus H^2(\bar{V}) \rightarrow H^2(U \cap \bar{V}) \rightarrow 0$$

Now $\boxed{h^0(U \cap \bar{V}) = h^1(U \cap \bar{V}) = 1, \quad h^2(U \cap \bar{V}) = 0}$

Also recall, that for an exact sequence of vector spaces $\sum (-1)^e a_e = 0$
 \uparrow
dim

Applying this remark to our sequence we immediately have:

notice this

$$\chi(M \# N) = \chi(U) - \chi(V) + 0 = 0, \text{ i.e.}$$

$$\boxed{\chi(M \# N) = \chi(U) + \chi(V)}$$

Now let $N = S^2$. Then $V = S^2 \setminus \{pt\} \approx \mathbb{R}^2$
 $\Rightarrow \chi(V) = 1$

Also $M \# S^2 \approx M \Rightarrow$

$$\chi(M) = \chi(U) + 1$$

Similarly (for generic N): $\chi(N) = \chi(V) + 1$

$$\Rightarrow \boxed{\chi(U) = \chi(M) - 1}, \quad \boxed{\chi(V) = \chi(N) - 1}$$

Therefore

$$\boxed{\chi(M \# N) = \chi(M) + \chi(N) - 2} \quad (*)$$

We now argue by induction: since $\Sigma_g = \Sigma_1 \# \dots \# \Sigma_1$
 (given the result for $g=1$ (already proven) $(g \geq 1)$ g copies

we have

$$\chi(\Sigma_g) = \chi(\Sigma_{g-1} \# \Sigma_1) = 2 - 2(g-1) - 2$$

$$= 2 - 2g \quad \boxed{\chi(\Sigma_g) = 2 - 2g}$$

* Remark. Notice that formula (*) is immediate in terms of triangulations: in performing the connected sum (see fig.) we "lose" 2 faces, 3 vertices and 3 edges, but the last two objects enter the Euler-Poincaré characteristic with opposite signs, and the assertion follows.



* Let M be a compact, connected orientable manifold, with no boundary

Then (e.g. via Poincaré Duality and $H^0(M) \cong \mathbb{R}$)

$$\boxed{H^n(M) \cong \mathbb{R}}$$

As a generator of the cohomology we may take a volume form ω (i.e. a never vanishing n -form, giving the orientation: choose it in such a way that

$$\int_M \omega > 0 \quad (\text{the integral is finite})$$

[Clearly, we may set the integral to 1 after multiplying by a suitable constant]

In fact, we show that $[\omega] \neq 0$.

By contradiction, if $[\omega]$ were equal to 0, then $\omega = d\alpha$, for $\alpha \in \Lambda^{n-1}(M)$. Moreover, by Stokes' Theorem:

$$0 < \int_M \omega \stackrel{=}{=} \int_M d\alpha \stackrel{\text{Stokes}}{=} \int_{\partial M} \alpha = 0,$$

yielding a contradiction.

* as a chain...

* Theorem: There are no symplectic structures on S^4

Recall that a symplectic manifold (M, ω) is a smooth manifold equipped with a closed, non degenerate 2-form ω (we already know then that $\dim M = \text{even}$)

Proof. Let ω be a symplectic form on S^4 : $d\omega = 0$

Then $[\omega] \in H^2(S^4) = \{0\}$ (recall...),

i.e. $\omega = d\alpha$, $\alpha \in \Lambda^1(S^4)$

Now

$$\begin{aligned}\omega \wedge \omega &= d\alpha \wedge \omega = d(\alpha \wedge \omega) + \alpha \wedge d\omega \\ &= d(\alpha \wedge \omega)\end{aligned}$$

But $\omega \wedge \omega$, in view of non degeneracy of ω , never vanishes, so it is a volume form, hence it cannot be exact, by the preceding observation.

$$\left(0 \neq \int_{S^4} \omega \wedge \omega = \int_{S^4} d(\alpha \wedge \omega) = \int_{\partial S^4} \alpha \wedge \omega = 0 \right)$$

* let us complete the cohomology of

$$\mathbb{R}^3 - \{(0,0,0)\}$$

clearly $\mathbb{R}^3 - \{(0,0,0)\} \underset{\text{Diffeo}}{\simeq} S^2 \times \mathbb{R}$

additional remarks:

$$H^0(\mathbb{R}^3 - \{0\}) \cong \mathbb{R}$$

connectedness

$$H^3(\mathbb{R}^3 - \{0\}) \cong 0$$

non compactness

$$H^2(\mathbb{R}^3 - \{0\}) \cong 0$$

is straightforward formal, since $\mathbb{R}^3 - \{0\}$

is simply connected, and one can use Poincaré lemma.

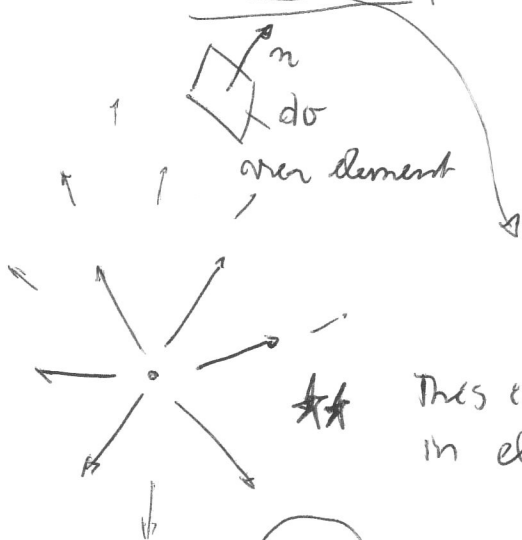
As a generator for the 2nd cohomology group one can

take

$$E = \frac{1}{4\pi} \frac{x dy dz + y dz dx + z dx dy}{(x^2 + y^2 + z^2)^{3/2}}$$

$$= \underline{E} \cdot \underline{n} d\sigma$$

= "flux element of an electric field generated by a point charge placed at the origin (charge value = $\frac{1}{4\pi}$ (normalization))"



$$\underline{E} = \frac{1}{4\pi} \frac{\underline{r}}{\|\underline{r}\|^3}$$

** This is the content of Gauss' Theorem in electrostatics



$$\iint_S \underline{E} \cdot \underline{n} d\sigma = 4\pi \cdot q$$

↑
charge

$$\iint_{S'} \underline{E} \cdot \underline{n} d\sigma = 0$$